

EUCAS2023
Bologna, Italy
3rd-7th September

Fe-based superconducting thin films and their potential for high field applications

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2023.09.05

Supported by

科研費 (16H04646, 20H02681)
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UNIVERSITY

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High-field transport properties: C. Tarantini, J. Hänisch
ASC NHMFL & Florida State University

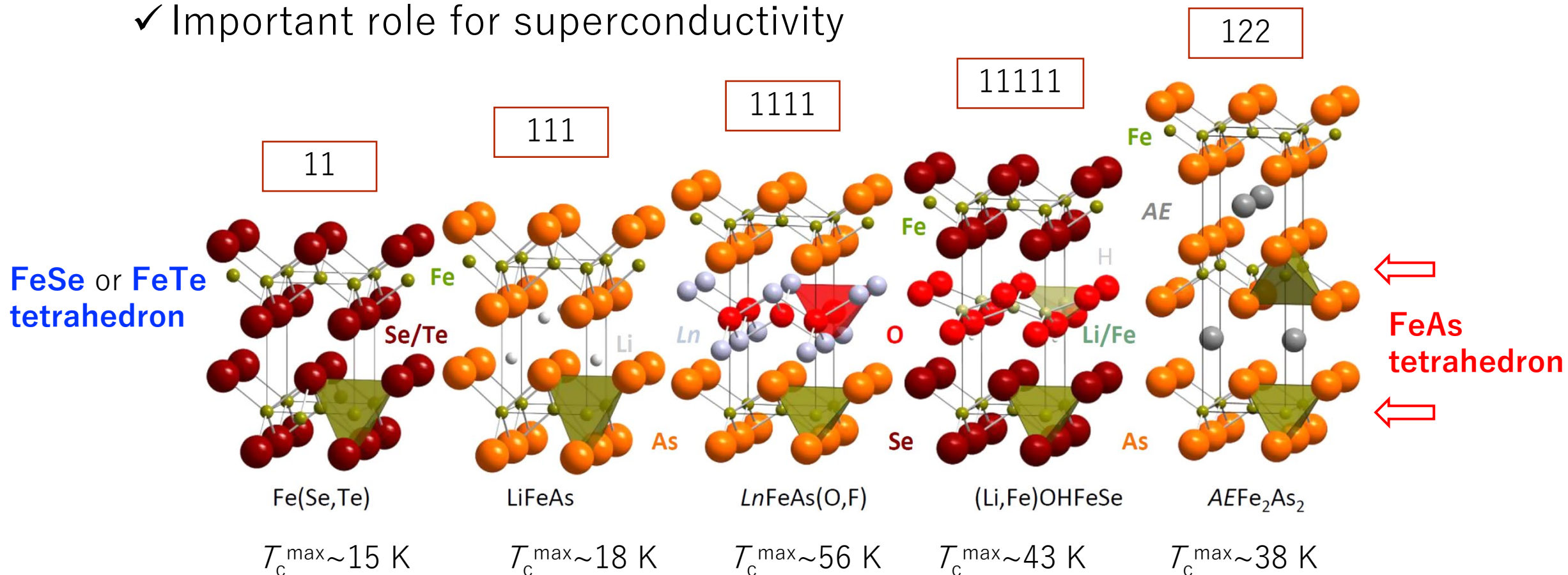


Inst. for Technical Physics, Karlsruhe Institute of Technology



FeAs or FeSe or FeTe tetrahedron

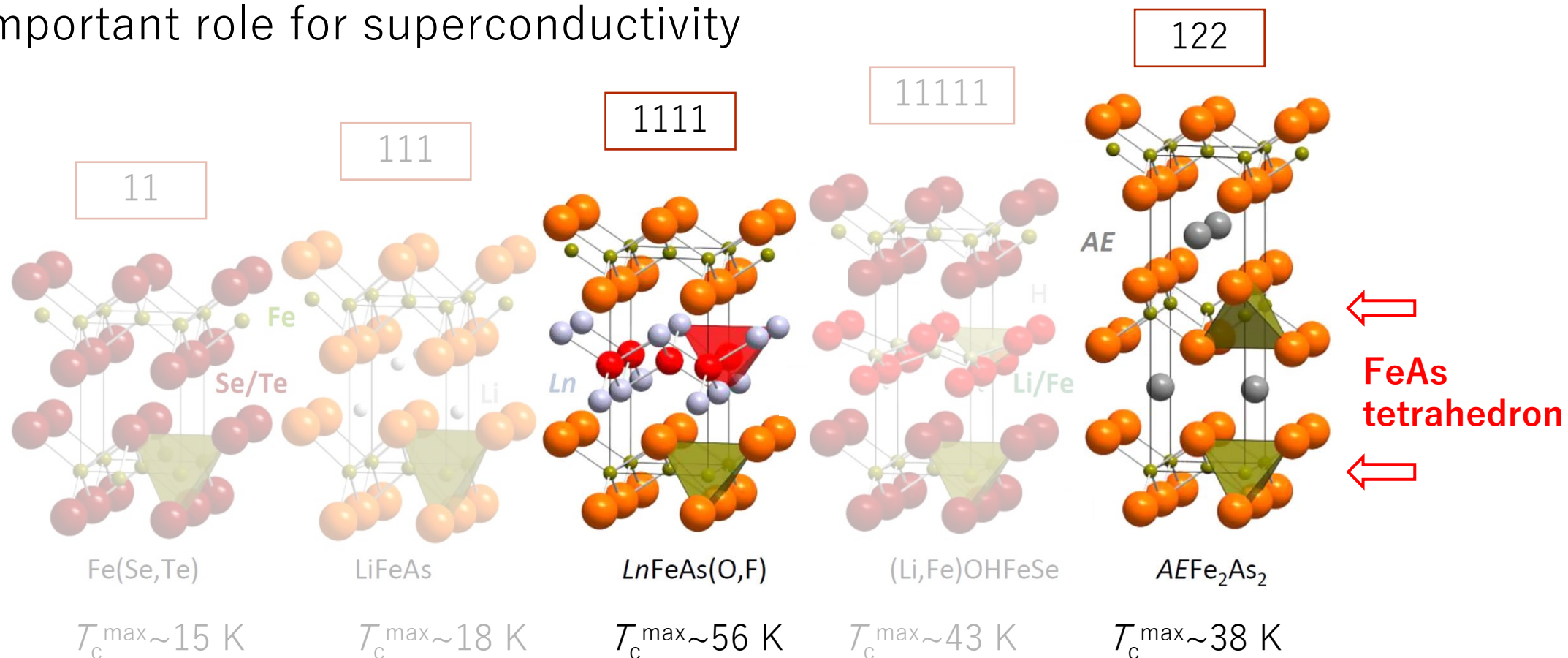
- ✓ Common structure
- ✓ Important role for superconductivity





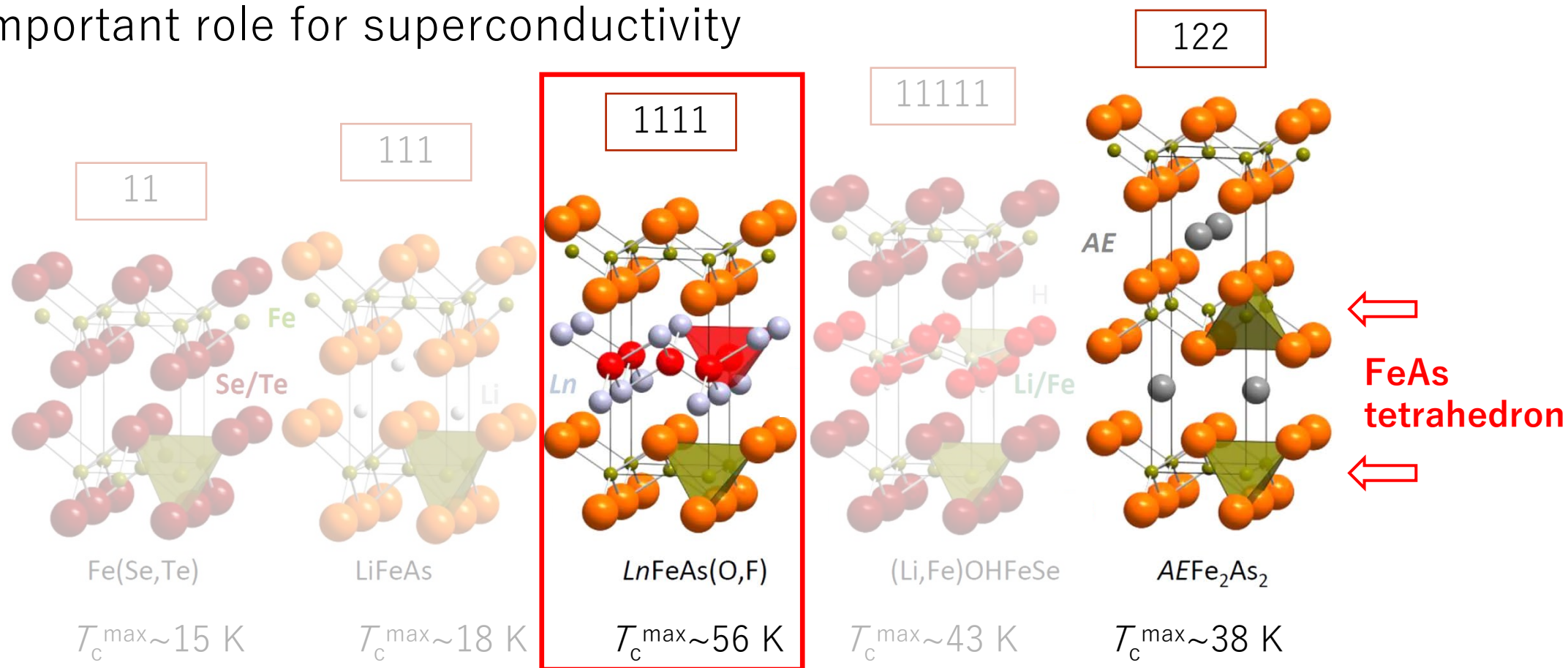
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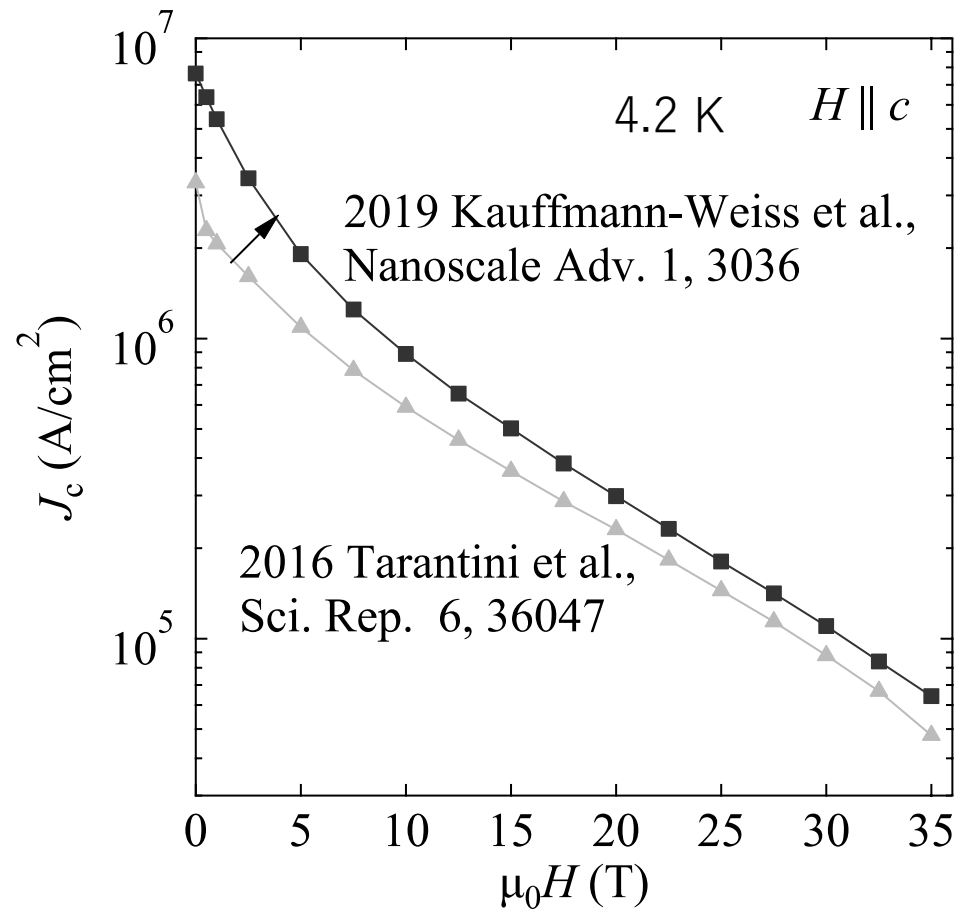


FeAs or FeSe or FeTe tetrahedron

- ✓ Common structure
- ✓ Important role for superconductivity

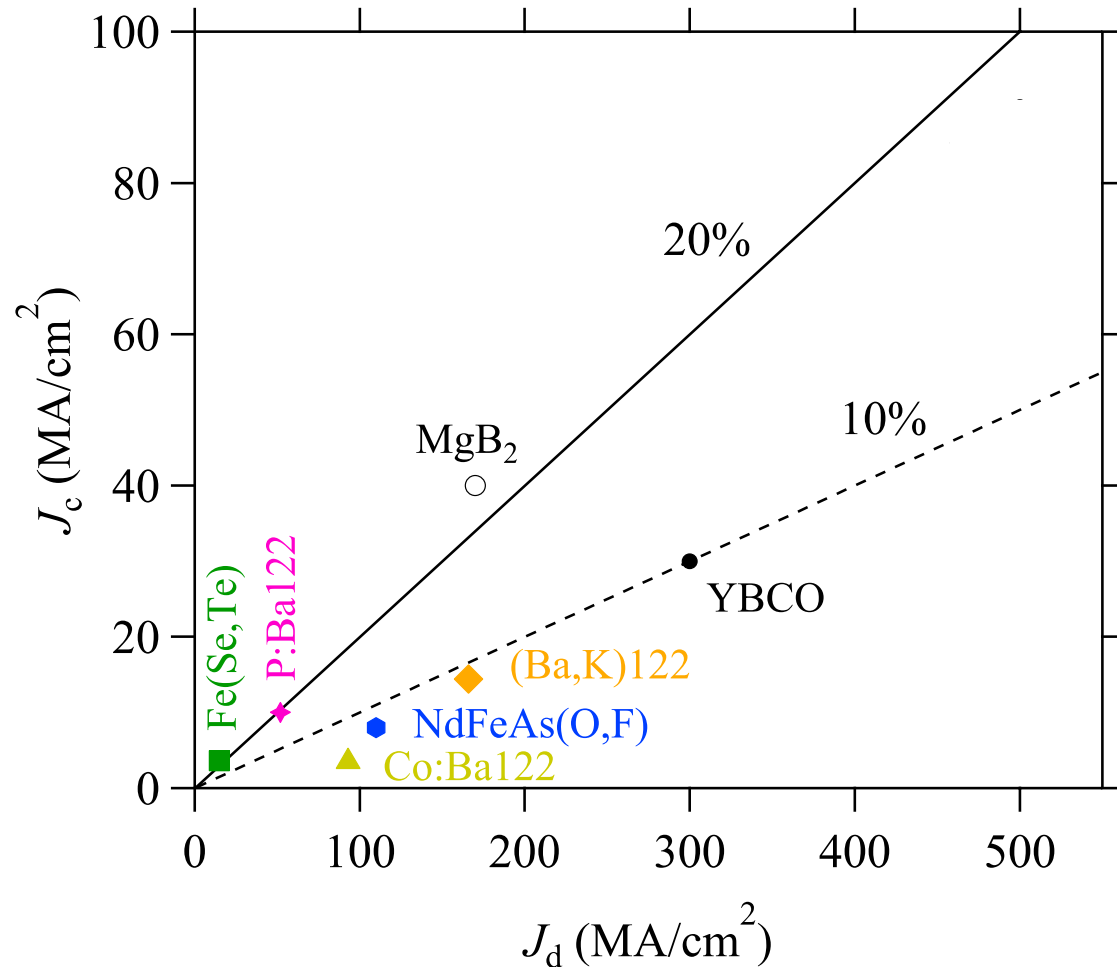


Progress of J_c - H for NdFeAs(O,F) (Nd-1111)



✓ $J_c \sim 8$ MA/cm² for NdFeAs(O,F)

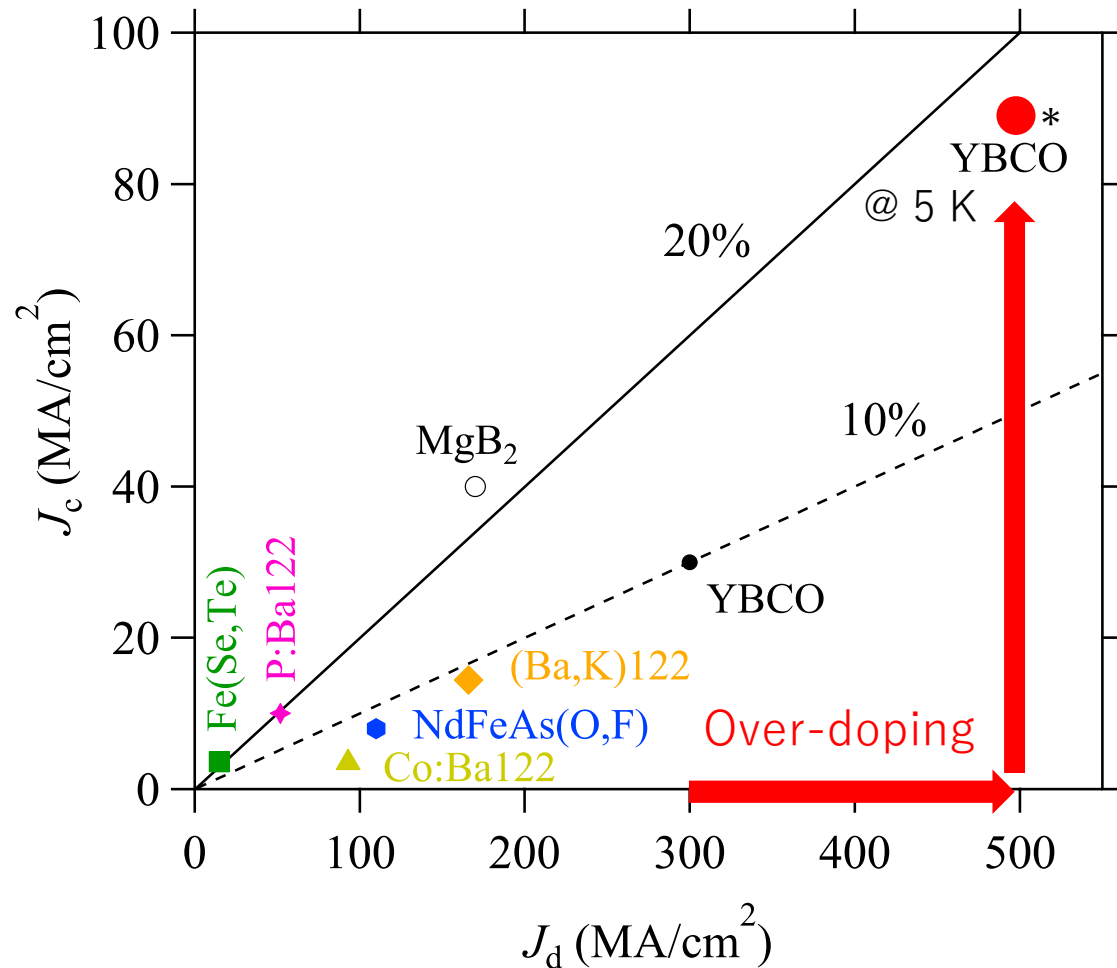
Measured J_c is around 10~20% of J_d



- ✓ $J_c \sim 8$ MA/cm² for NdFeAs(O,F)
- ✓ Higher the J_d , higher the J_c
- ✓ Depairing current density J_d
(theoretical upper limit)

$$J_d(T) = \frac{\phi_0}{3\sqrt{3}\pi\mu_0\lambda^2(T)\xi(T)}$$

Over-doping increases the condensation energy



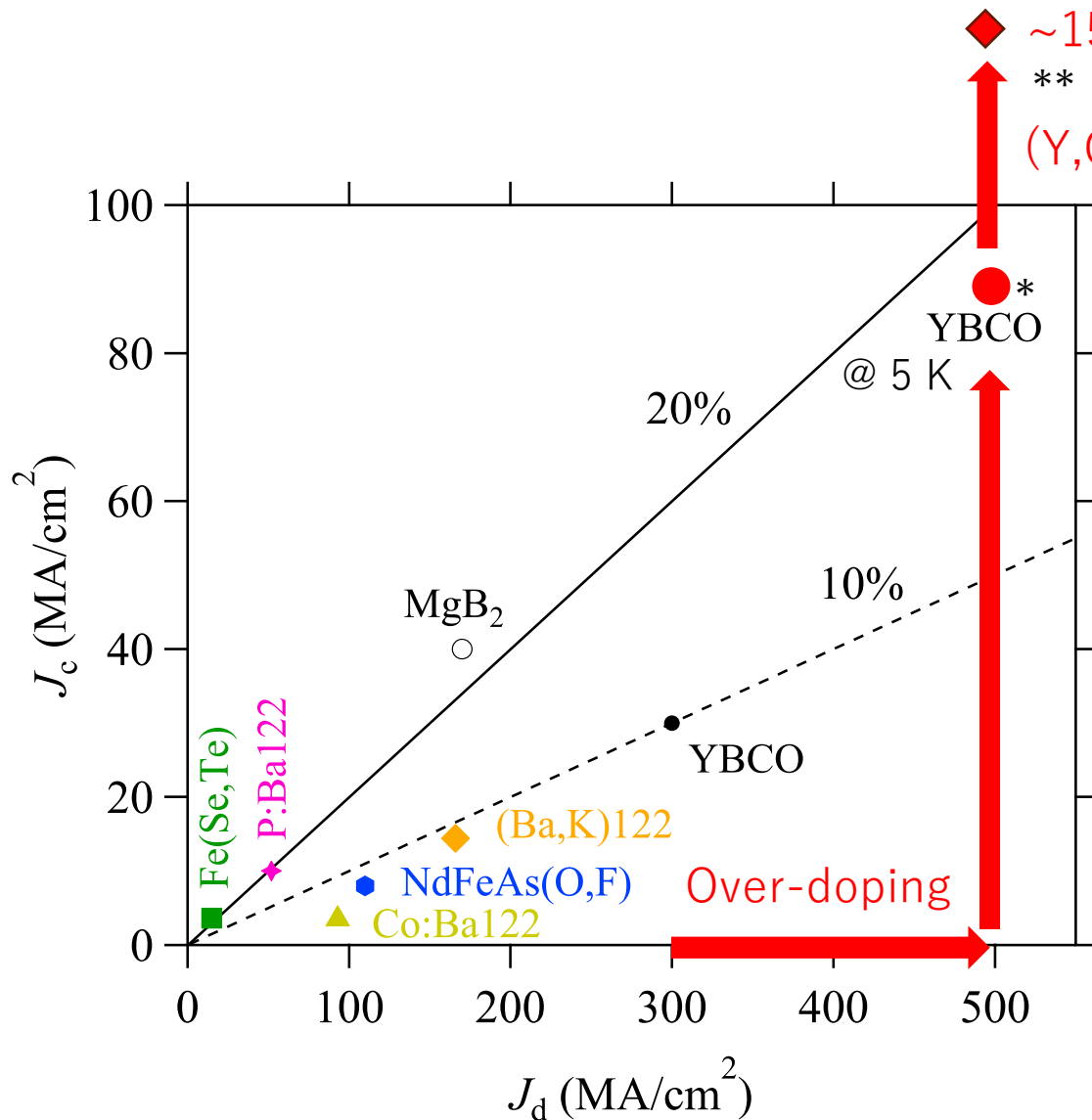
- ✓ $J_c \sim 8$ MA/cm² for NdFeAs(O,F)
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- ✓ Over-doping increase J_d , and hence J_c

* A. Stangl *et al.*, *Sci. Rep.* **11**, 8176 (2021).

Combining over-doping and microstructural modification



- ✓ $J_c \sim 8 \text{ MA/cm}^2$ for NdFeAs(O,F)
- ✓ Higher the J_d , higher the J_c
- ✓ Depairing current density J_d
(theoretical upper limit)

$$J_d(T) = \frac{\phi_0}{3\sqrt{3}\pi\mu_0\lambda^2(T)\xi(T)}$$

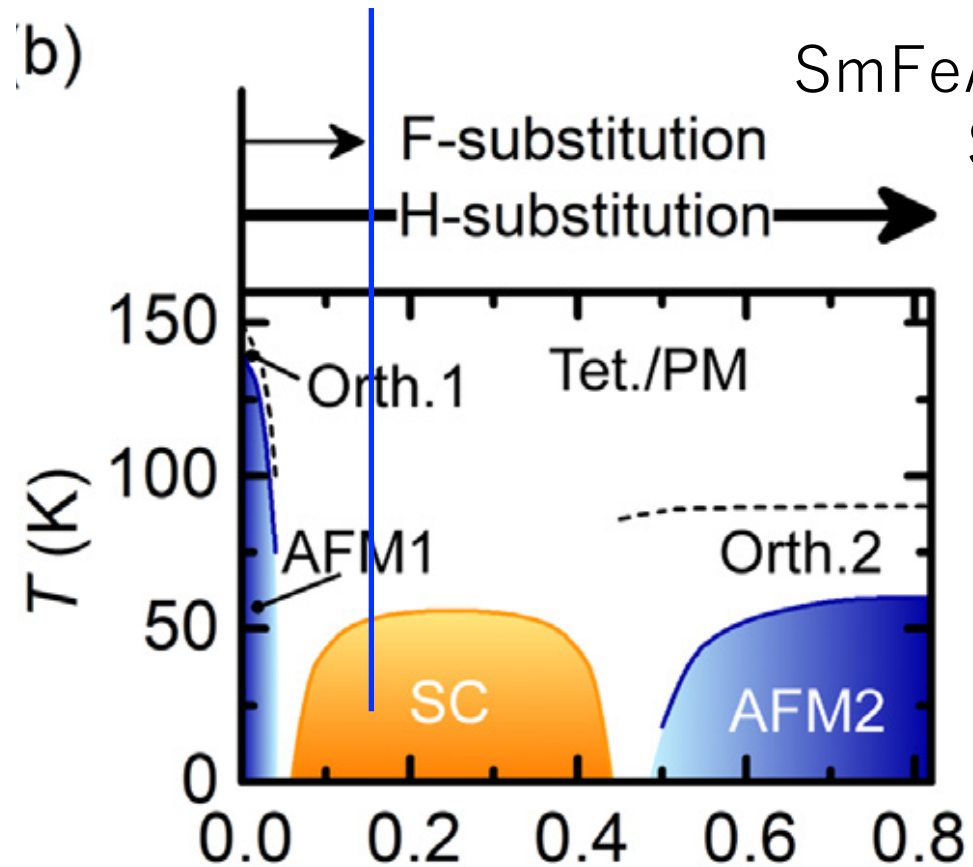
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* A. Stangl *et al.*, *Sci. Rep.* **11**, 8176 (2021).

** M. Miura *et al.*, *NPG Asia Mater.* **14**, 85 (2022).

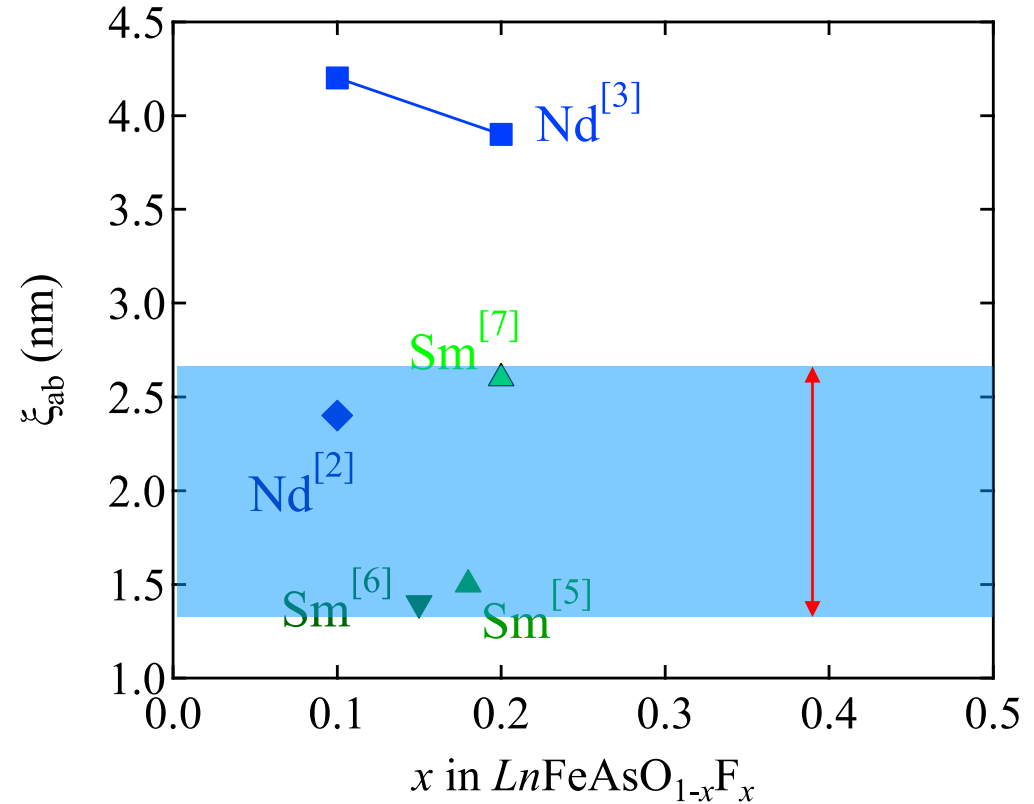
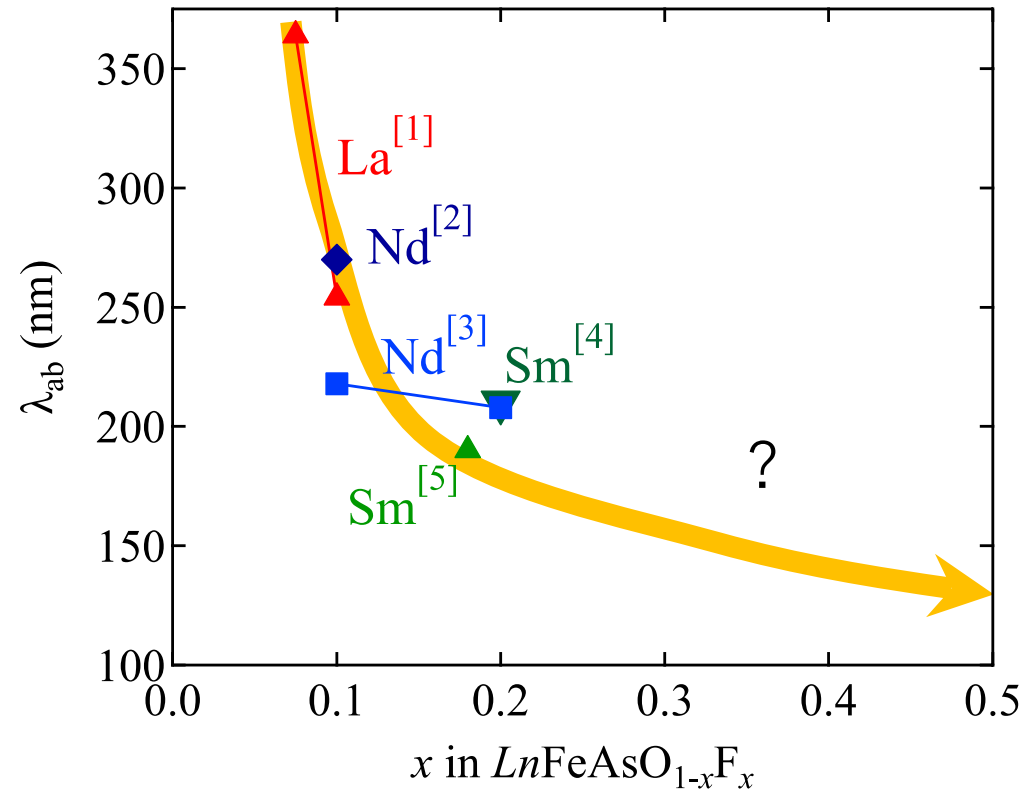
$LnFeAsO$ (Ln : lanthanoide)

$O^{2-} \rightarrow F^-$ or $H^- + e^-$ (electron doping)



- Substitution level is limited **up to ~0.2** (For SmFeAsO_{1-x}F_x)
- For H, the substitution level is increased **up to ~0.8**
- Heavily electron doped film can be obtained
- T_c keeps constant around 50 K up to $x=0.4$

$LnFeAsO$: penetration depth and coherence length



- λ of $LnFeAsO$ may be decreased with electron doping
- No clear relation between ξ and electron doping

[1] H. Luetkens *et al.*, *Phys. Rev. Lett.* **101**, 097009 (2008).

[2] J. Kacmarcik *et al.*, *Phys. Rev. B* **80**, 014515 (2009).

[3] A. Adamski *et al.*, *Phys. Rev. B* **96**, 100503 (2017).

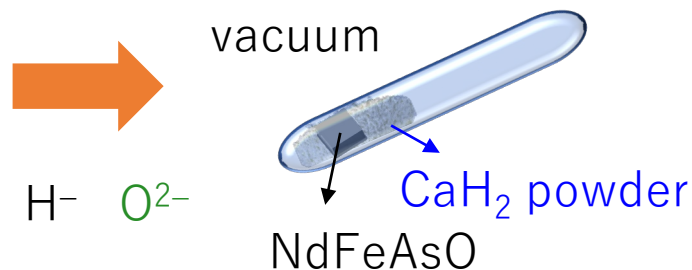
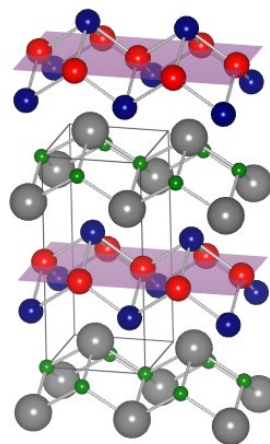
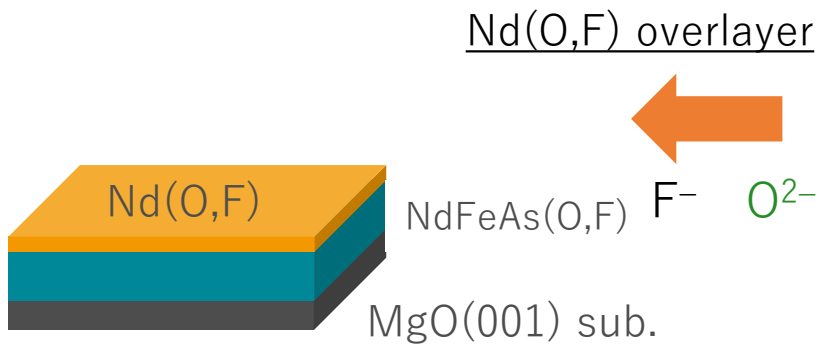
[4] S. Weyeneth *et al.*, *J. Supercond. Nov. Magn.* **22**, 325 (2009).

[5] A. J. Drew *et al.*, *Phys. Rev. Lett.* **101**, 097010 (2008).

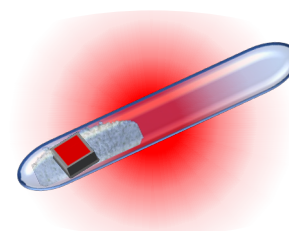
[6] U. Welp *et al.*, *Phys. Rev. B* **83**, 100513 (2011).

[7] H-S. Lee *et al.*, *Phys. Rev. B* **80**, 144512 (2009).

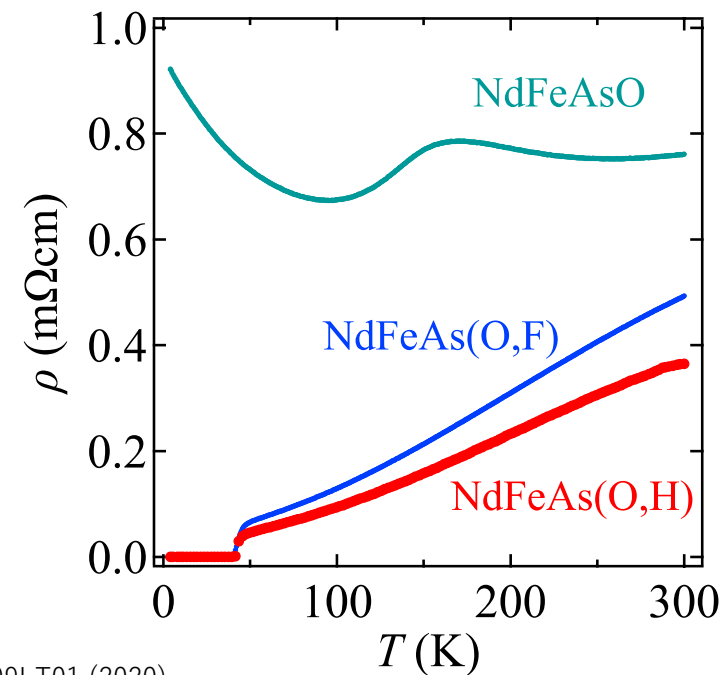
Growth of F- and H-doped NdFeAsO thin films



heat treatment

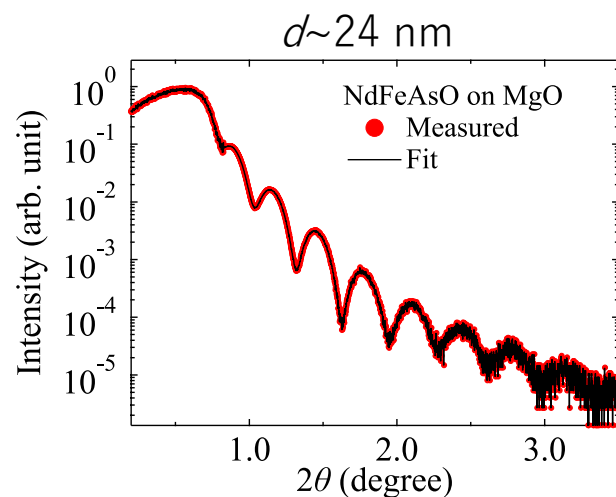
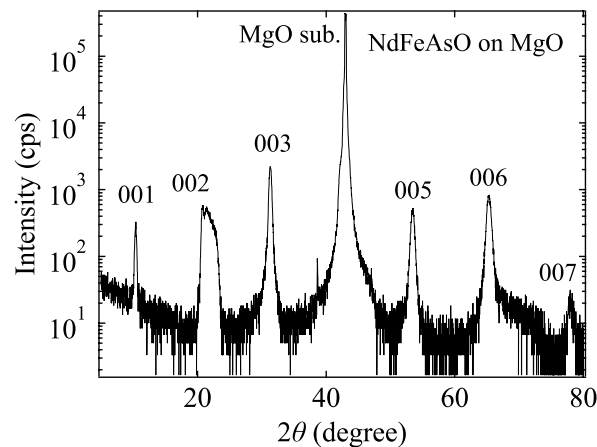


400~500°C
up to 72 h



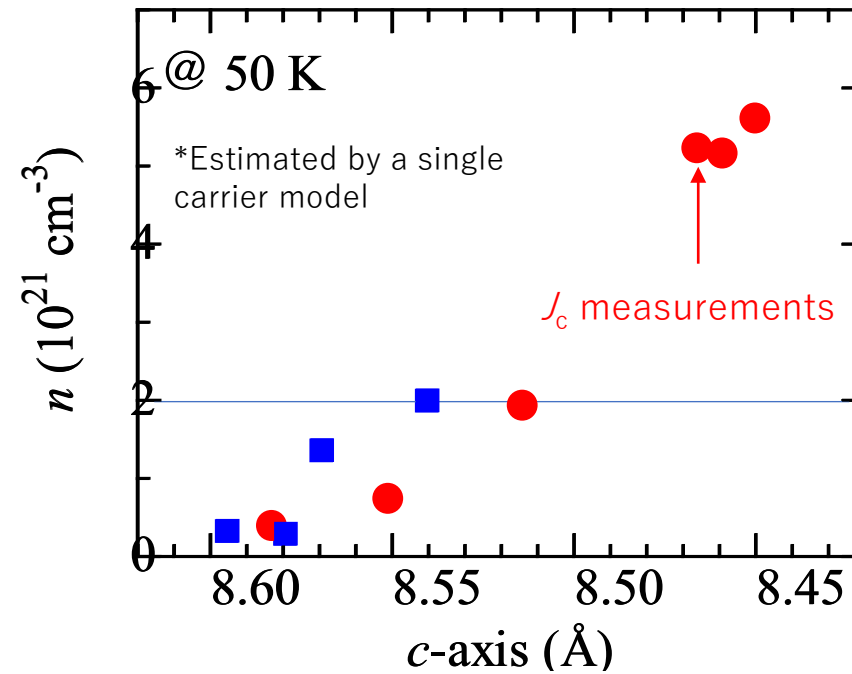
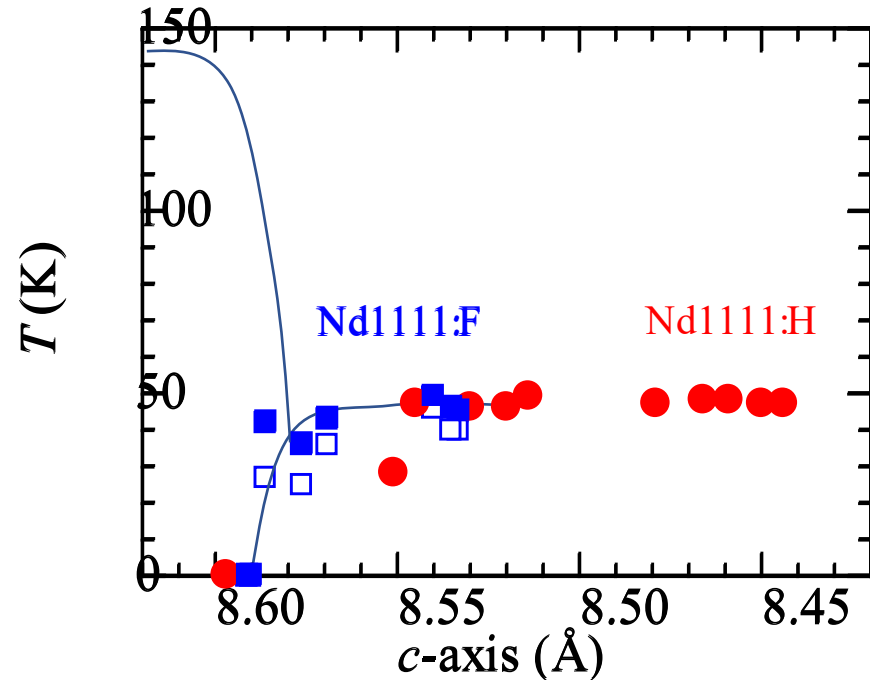
T. Kawaguchi *et al.*, *Appl. Phys. Express* **4**, 083102 (2011).

Phase pure & *c*-axis oriented



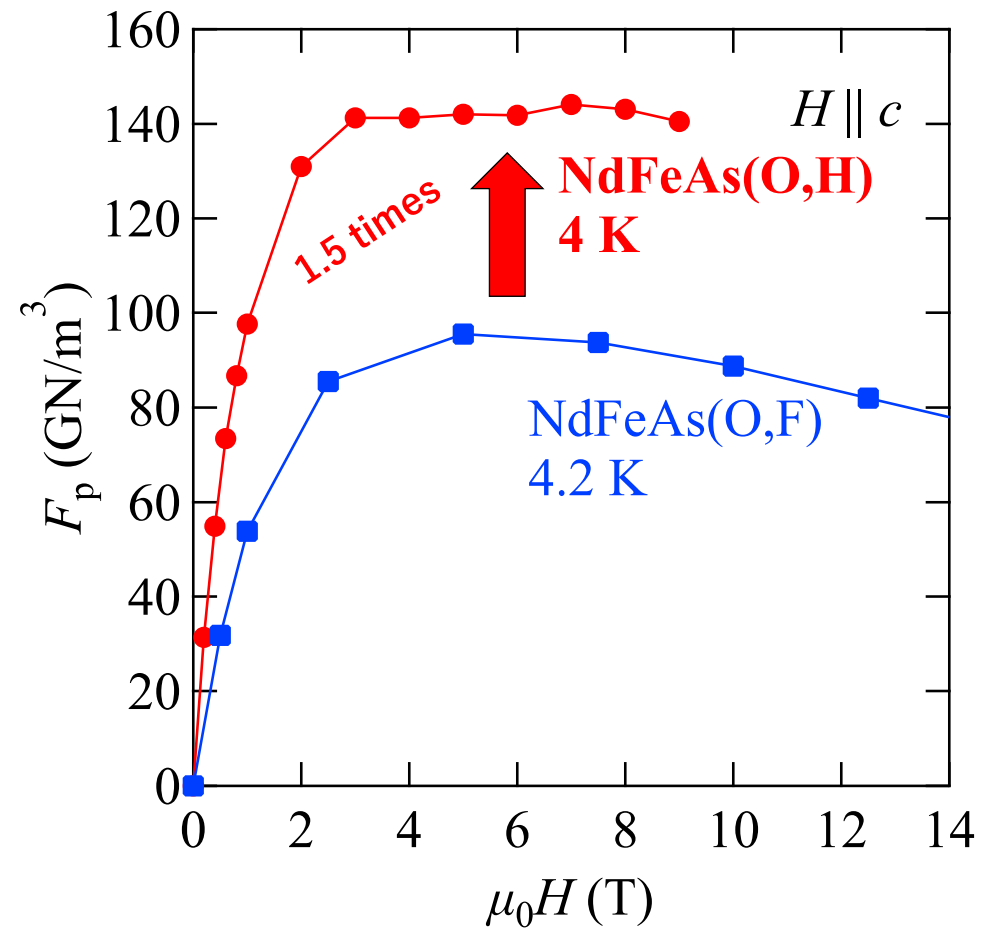
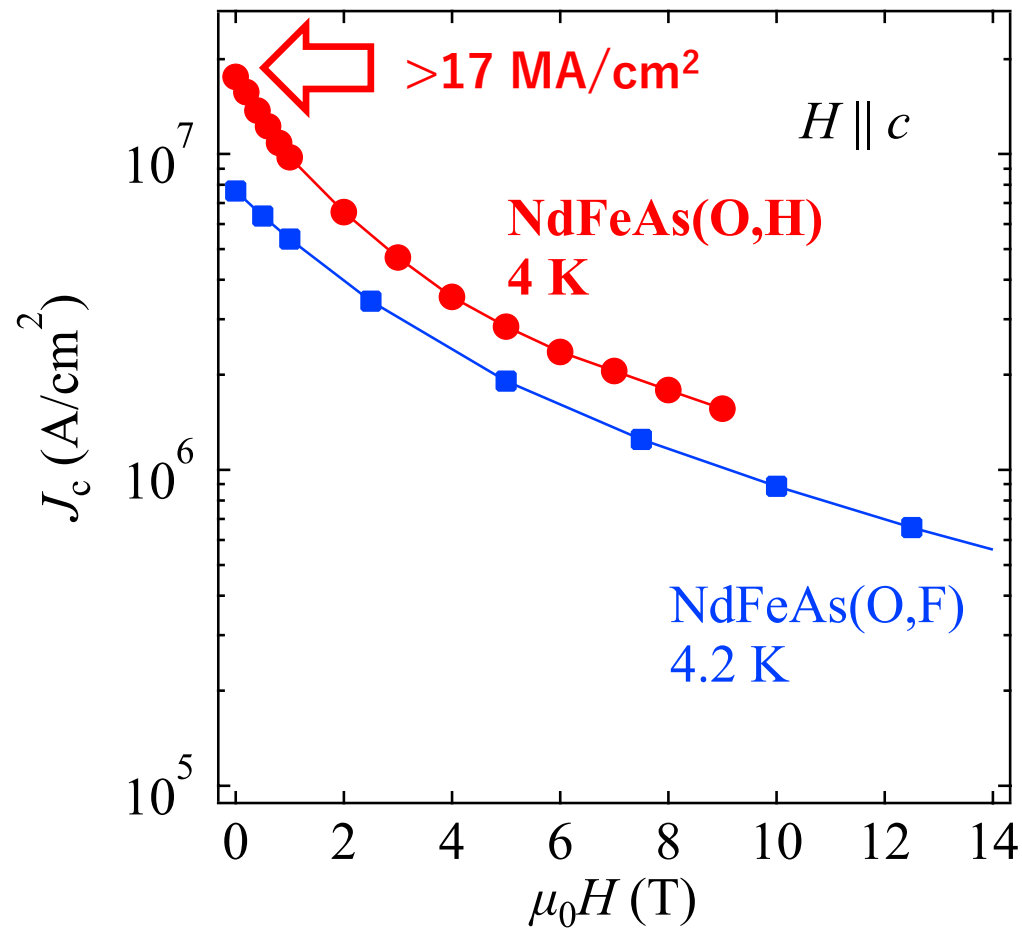
K. Kondo, K. Iida *et al.*, *SuST* **33**, 09LT01 (2020).

Comparison between H- and F-doping: T_c and carrier density



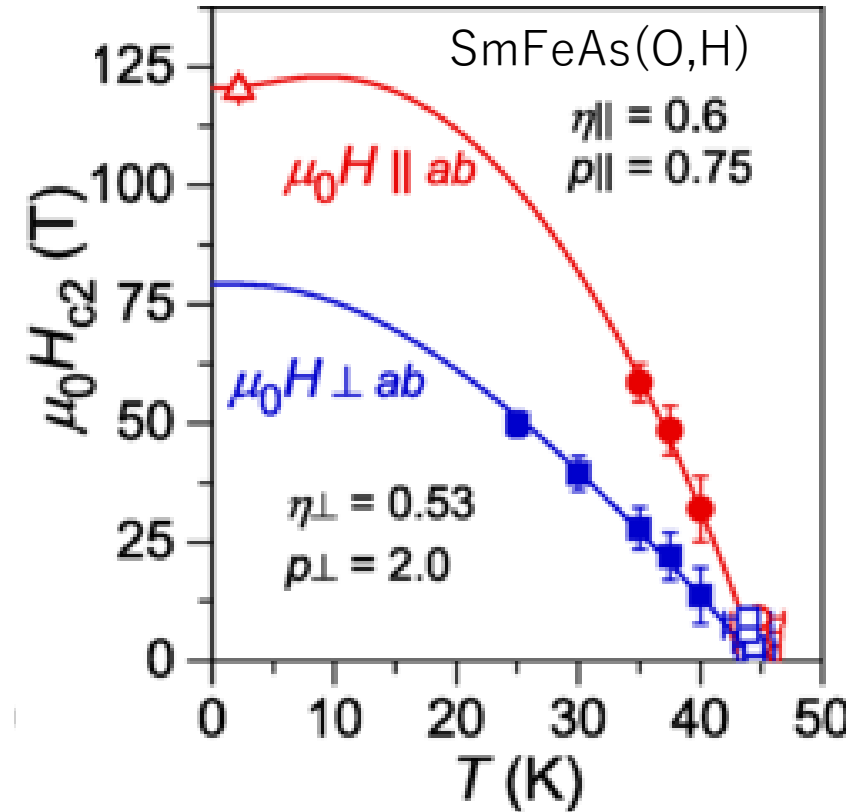
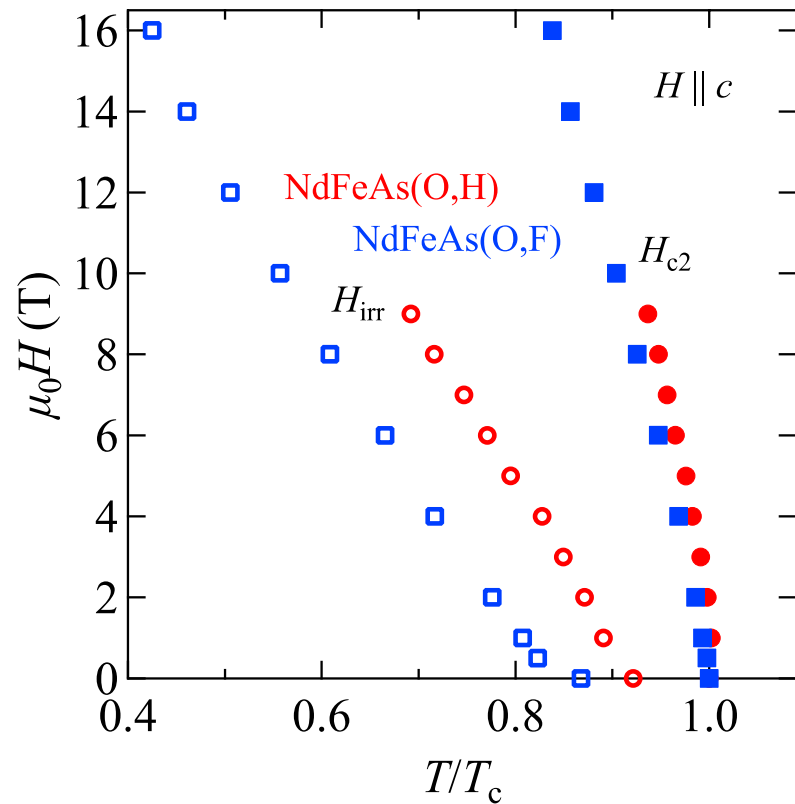
- A constant T_c of ~ 50 K up to $c \sim 8.54$ \AA , corresponding to carrier density $n \sim 2 \times 10^{21} / \text{cm}^3$
- c varied in a wide range, 8.44 \AA and 8.55 \AA with $T_c \sim 50$ K [NdFeAs(O,H)]
- The maximum n was around $6 \times 10^{21} / \text{cm}^3$ [almost 3 times higher than NdFeAs(O,F)]

J_c - H for NdFeAs(O,H) is higher than that for NdFeAs(O,F)



- $J_c(4 \text{ K}, 0 \text{ T}) > 17 \text{ MA}/\text{cm}^2$, which is twice larger than NdFeAs(O,F) c.f.) $20 \text{ MA}/\text{cm}^2$ for the irradiated SmFeAs(O,F) single crystal [1]
- F_p of NdFeAs(O,H) is 1.5 times higher than that of NdFeAs(O,F)

NdFeAs(O,H) shows higher H_{irr} than NdFeAs(O,F)

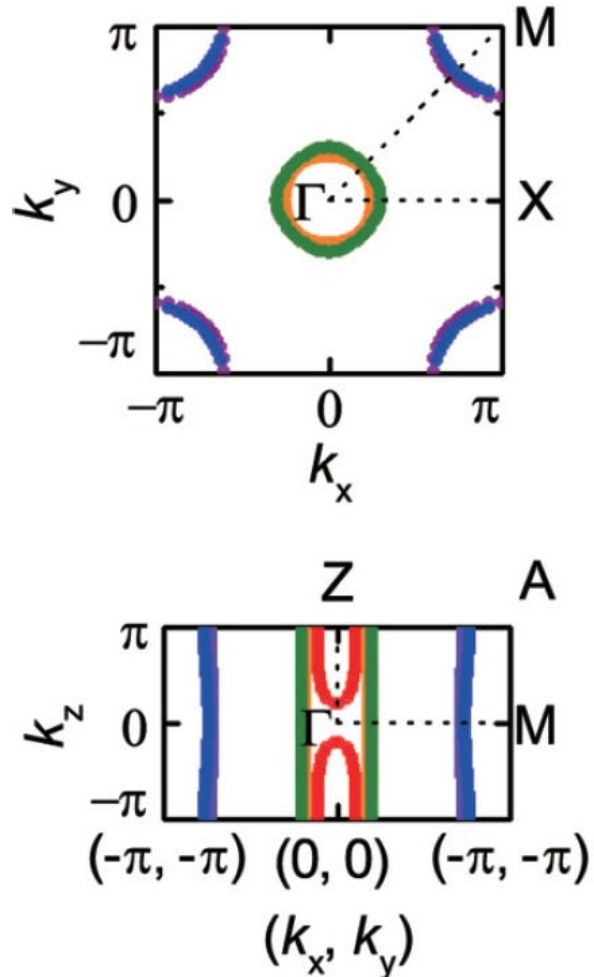


Hanzawa *et al.*, *PRM* **6**, L1118 (2022).

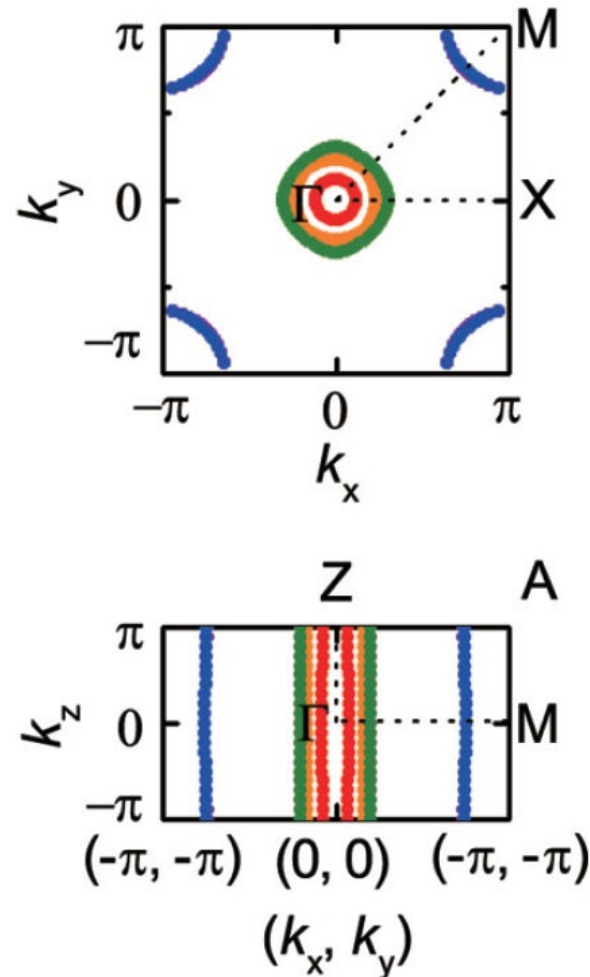
- Upper critical field H_{c2} of NdFeAs(O,H) is almost comparable to that of NdFeAs(O,F)
- Irreversibility field H_{irr} of NdFeAs(O,H) is higher than that of NdFeAs(O,F)

-> Due to the decrease in the anisotropy $H_{irr} \propto \frac{H_{c2}}{\gamma^2}$ for $H \parallel c$

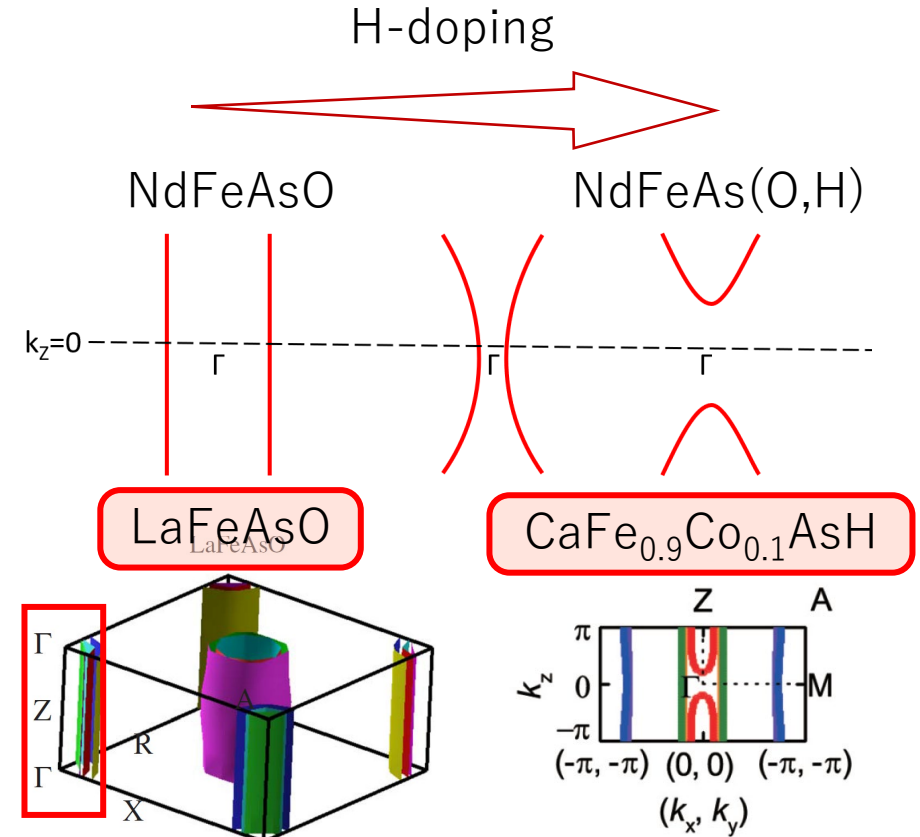
Fully substituted H and F of 1111



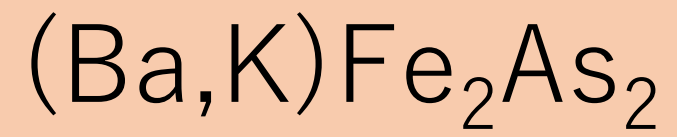
Y. Muraba *et al.*, *PRB* **89**, 094501 (2014).



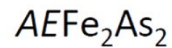
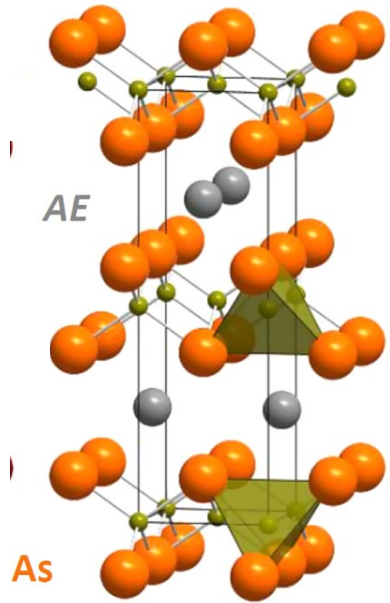
Evolution of band structure with H



V. Vildosola *et al.*, *PRB* **78**, 064518 (2008).

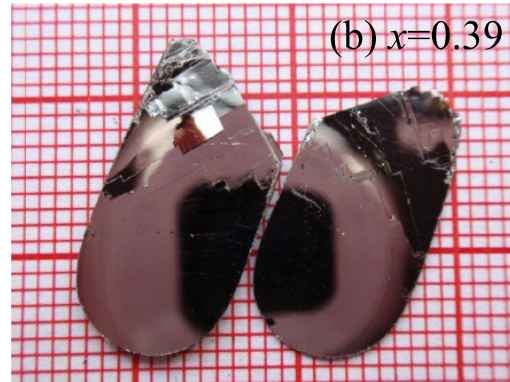


K-doped AFe_2As_2 : Only epitaxial thin film has not been realised



- ✓ $T_c^{\max} \sim 38$ K
- ✓ $J_d \sim 170$ MA/cm²
- ✓ $\gamma \sim 1 - 2$

Single crystals

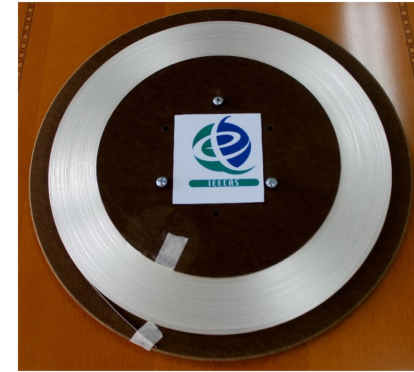


Y. Liu *et al.*, *PRB* **89**, 13504 (2014).

Polycrystalline bulk

J. Weiss *et al.*, *SuST* **28**, 112001 (2015).

Wires

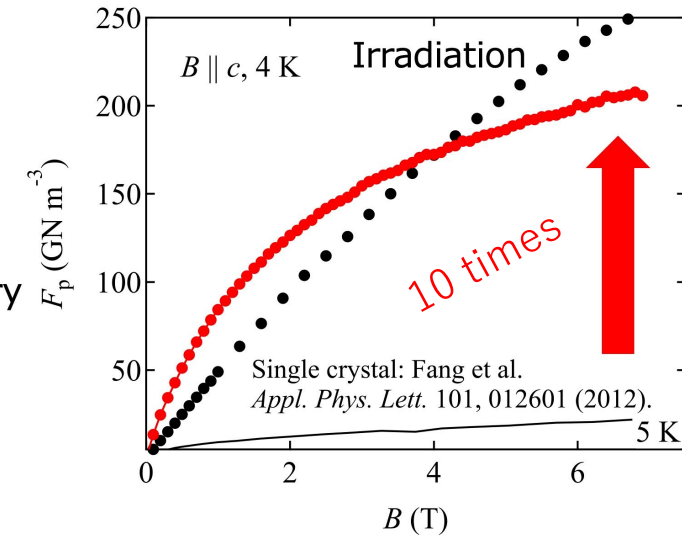
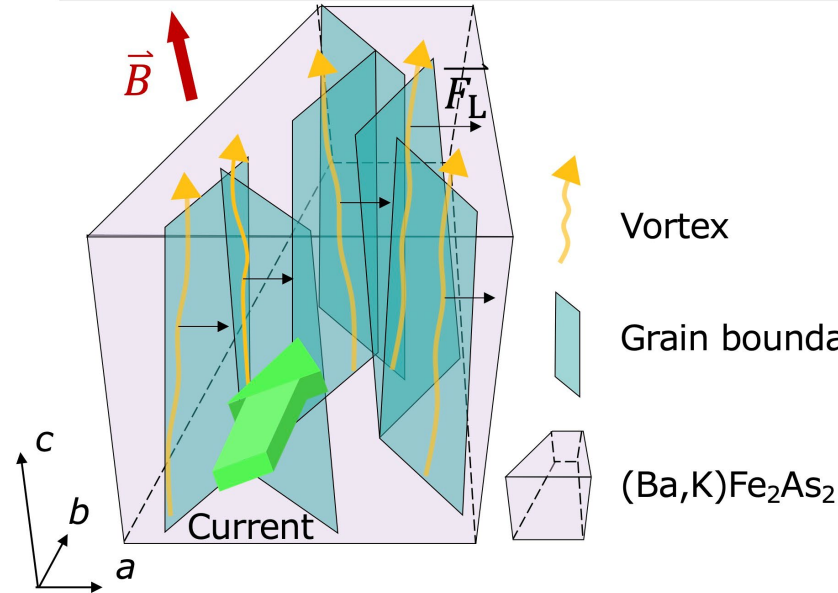
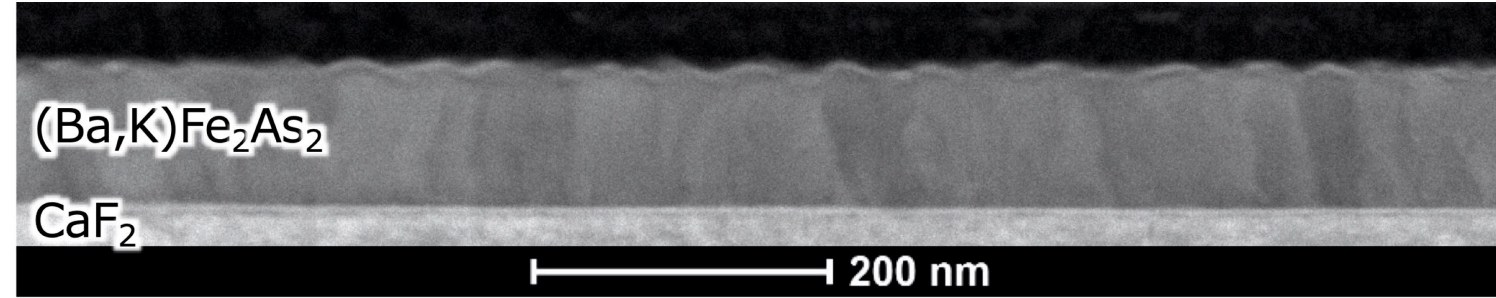
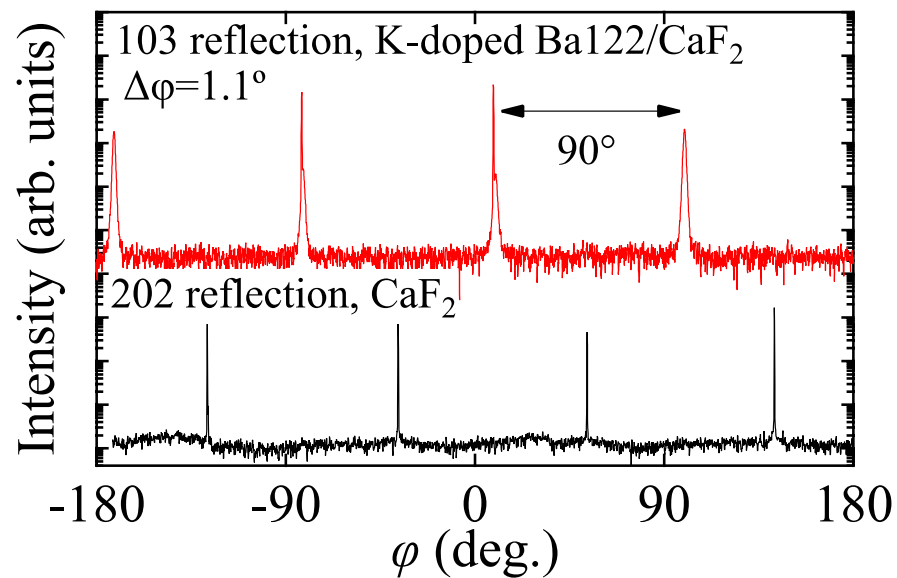
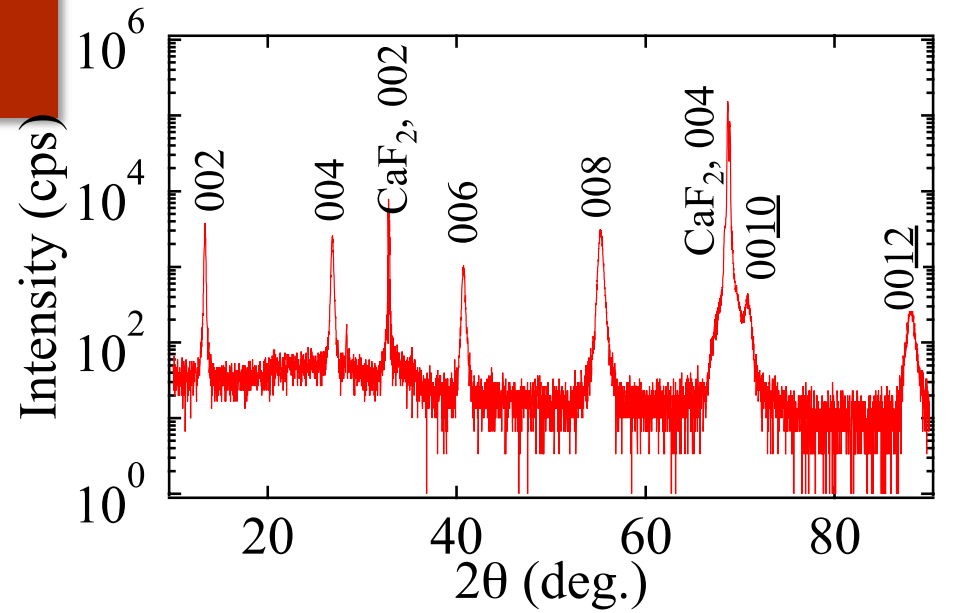


H. Hosono *et al.*, *Mater. Today* **21**, 278 (2018).



- ✓ High $T_c \sim 38$ K, high $J_d \sim 170$ MA/cm² & low electromagnetic anisotropy
- ✓ The most promising material for applications (e.g. PIT wires and bulk magnets)
- ✓ All material forms except for epitaxial thin films have been available

K-doped Ba122 epitaxial thin film on CaF₂ sub.



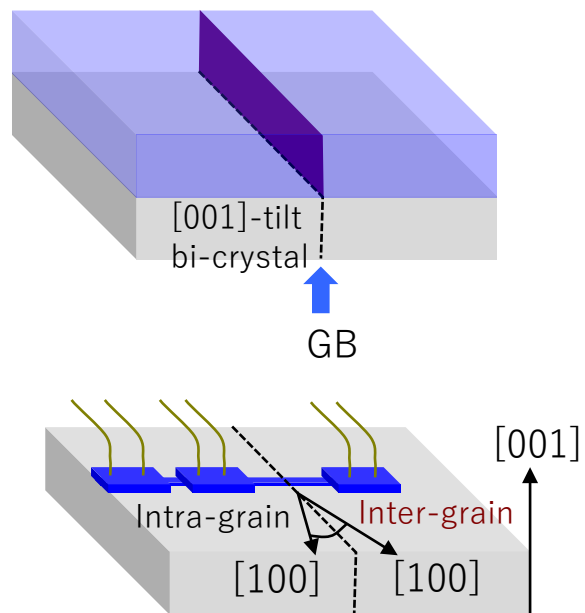
Fluoride sub. & low growth temperature
(~400° C) are the key to epitaxial growth

D. Qin *et al.*, *Phys. Rev. Mater.* **5**, 014801 (2021).

K. Iida *et al.*, *NPG Asia Materials* **13**, 68 (2021).



Toward bi-crystal K-doped Ba122 experiments

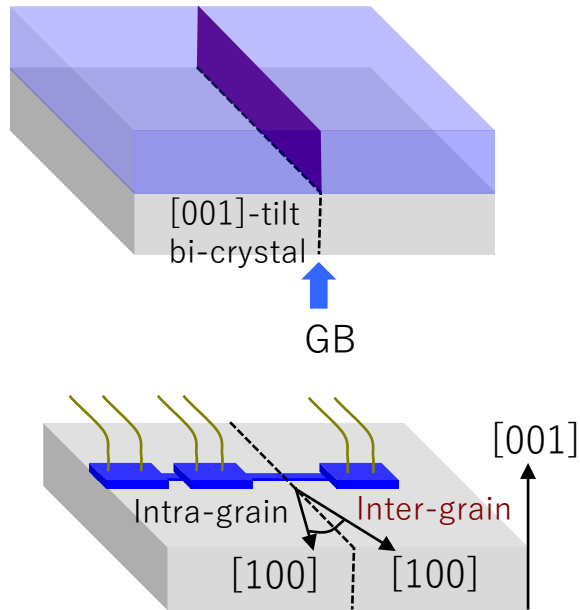


Substrates	Bi-crystal
CaF ₂ (001)	Not available
MgO (001)	Available
SrTiO ₃ (001)	Available

Substrates	Bi-crystal
LSAT ¹⁾ (001)	Available
Al ₂ O ₃ (0001)	Available

¹⁾ La_{0.3}Sr_{0.7}Al_{0.65}Ta_{0.35}O₃

Toward bi-crystal K-doped Ba122 experiments

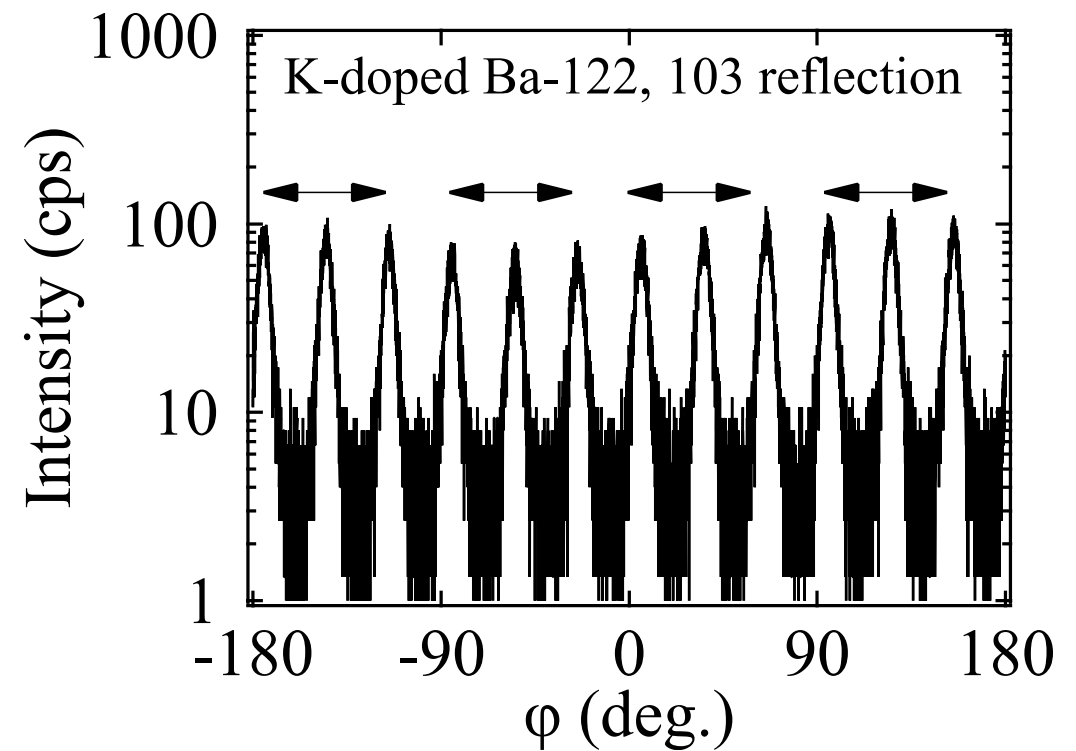
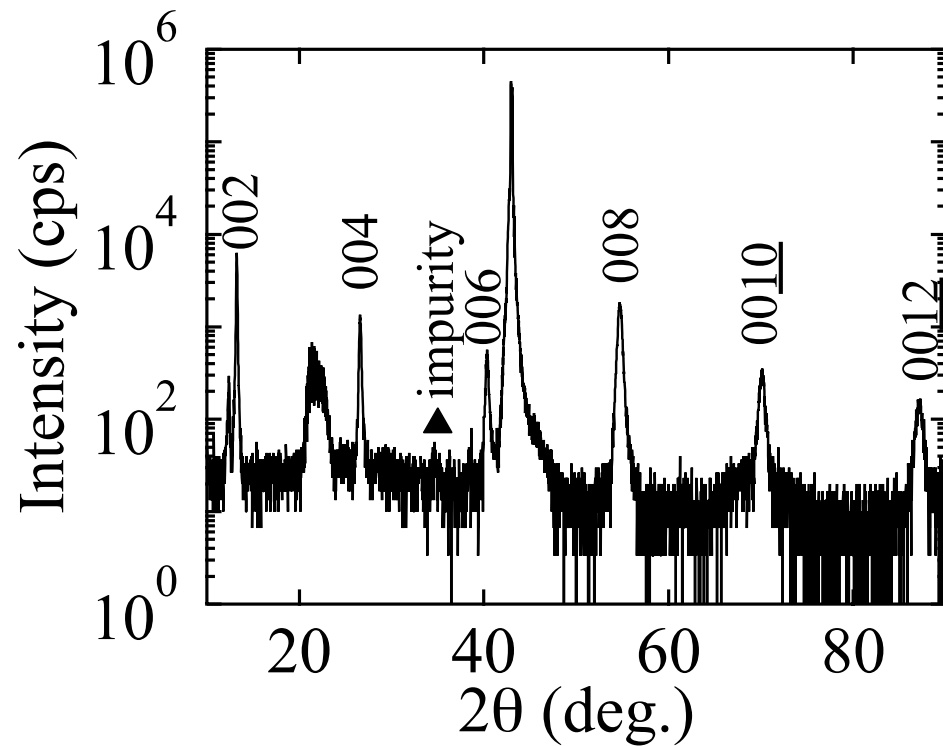


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SrTiO ₃ (001)	Available

Substrates	Bi-crystal
LSAT (001)	Available
Al ₂ O ₃ (0001)	Available



- ≈ Oxygen is released during deposition (SrTiO₃)
- ≈ Reaction layer is present at the interface (LSAT)
- ☉ Ba122 grows epitaxially on MgO

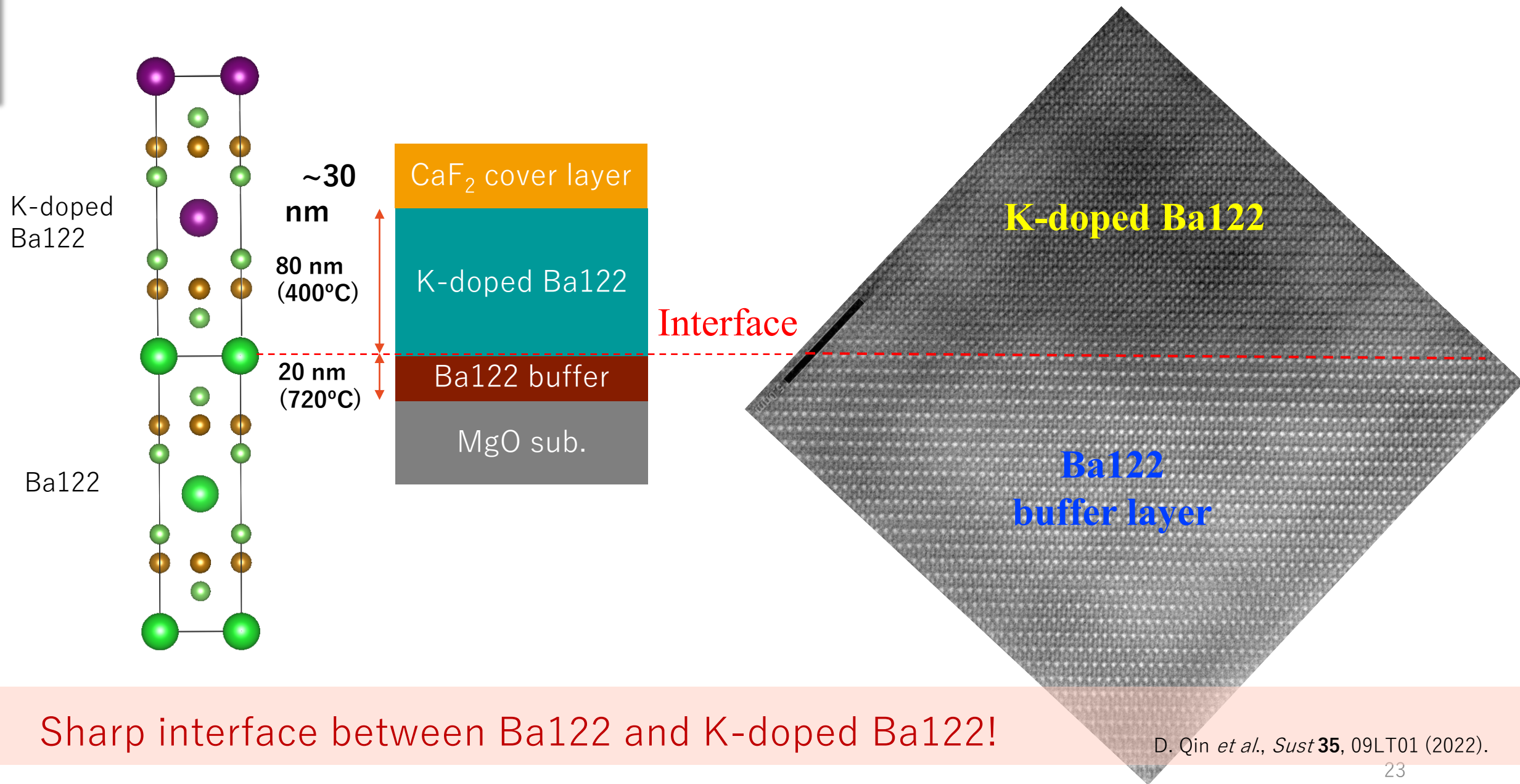


- $(\text{Ba},\text{K})\text{Fe}_2\text{As}_2$ was not epitaxially grown \rightarrow due to the low growth temperature (cf. 800°C for Co-^[1] and P-doped^[2] Ba122)
- 30° rotated grains were present

[1] T. Katase *et al.*, *SuST* **25**, 084015 (2012).
[2] S. Adachi *et al.*, *SuST* **25**, 105015 (2012).

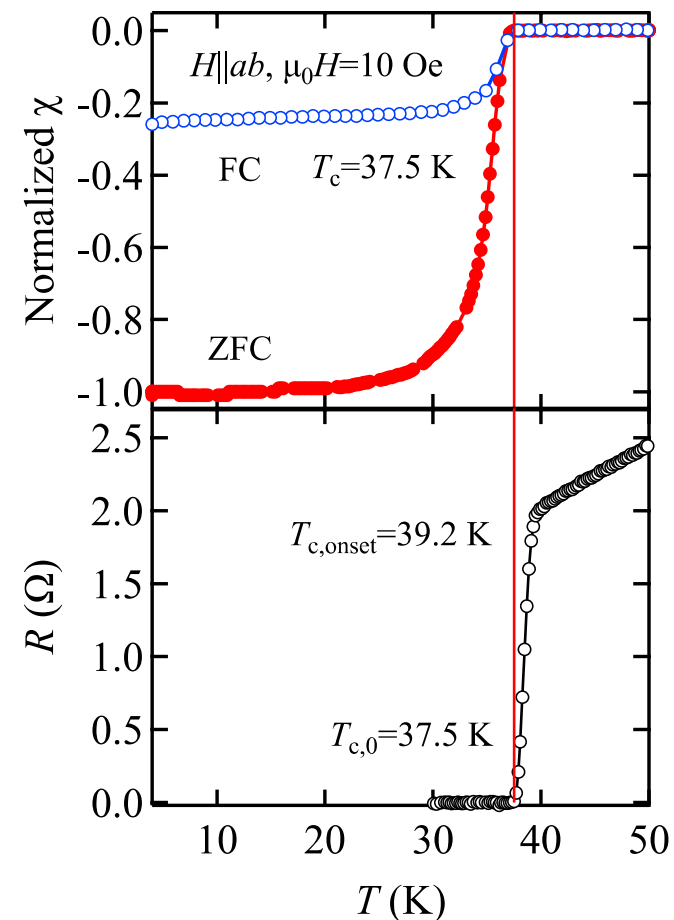
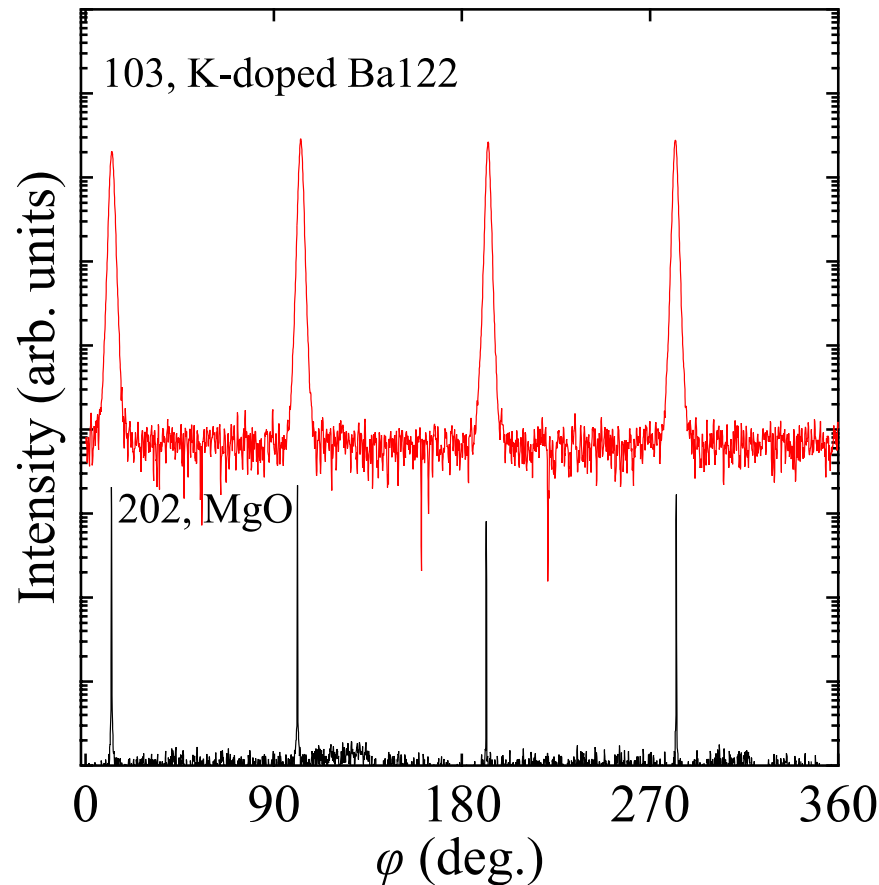
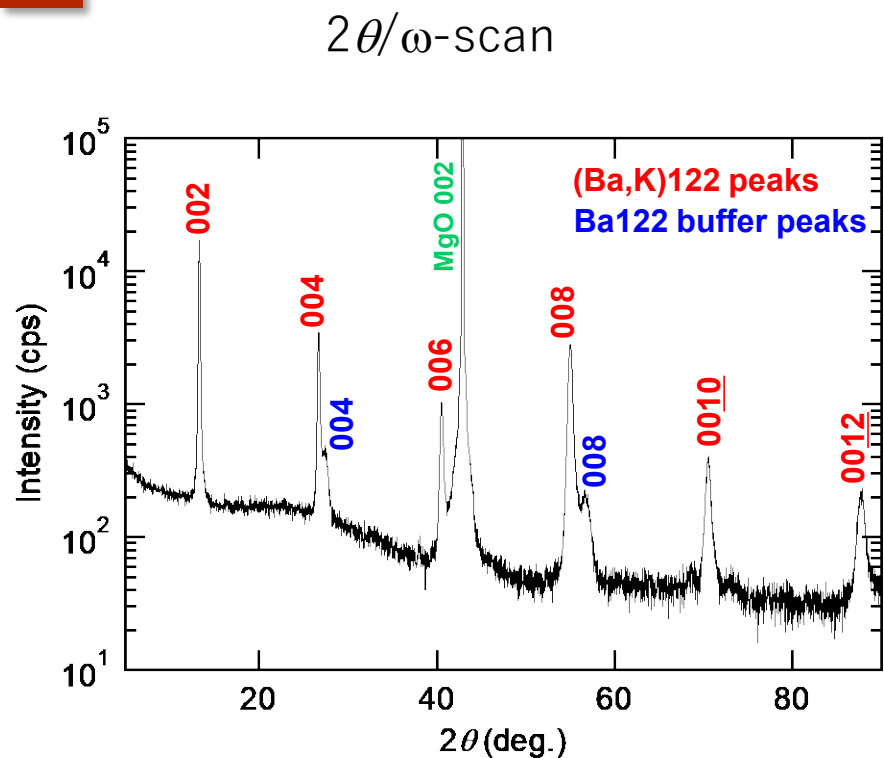


Strategy for K-doped Ba122 on MgO



Sharp interface between Ba122 and K-doped Ba122!

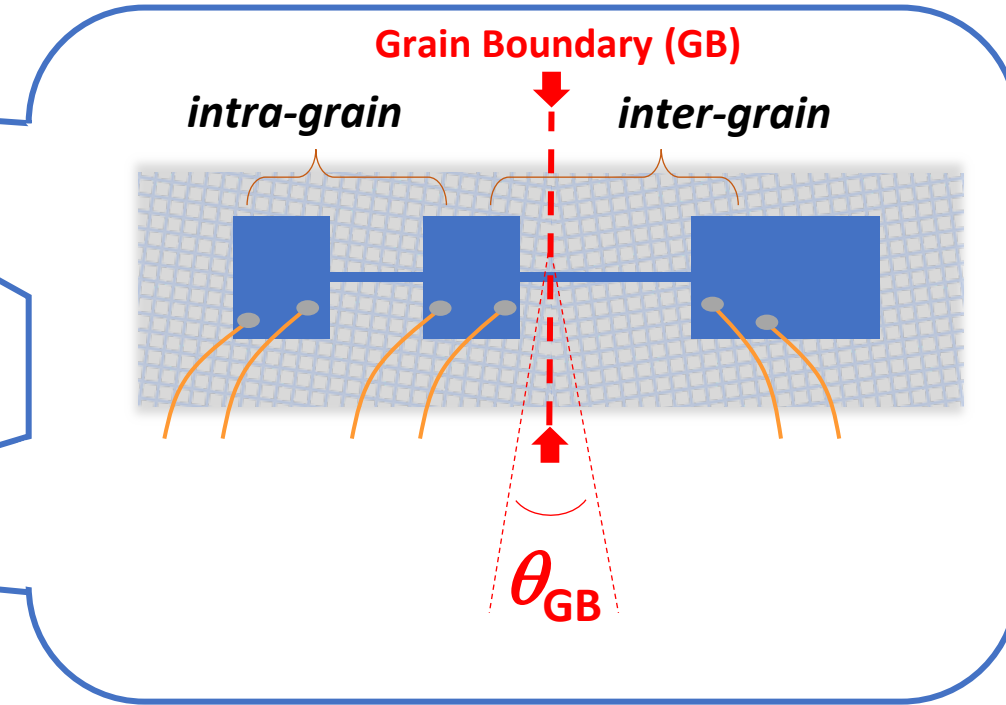
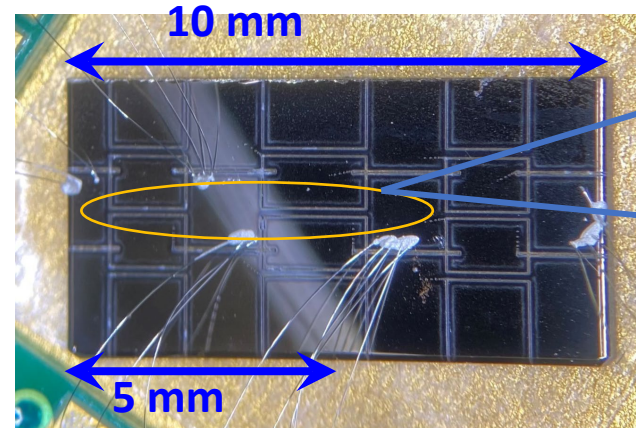
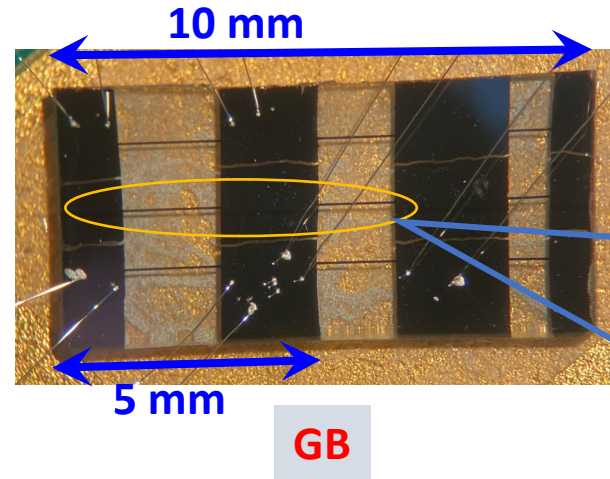
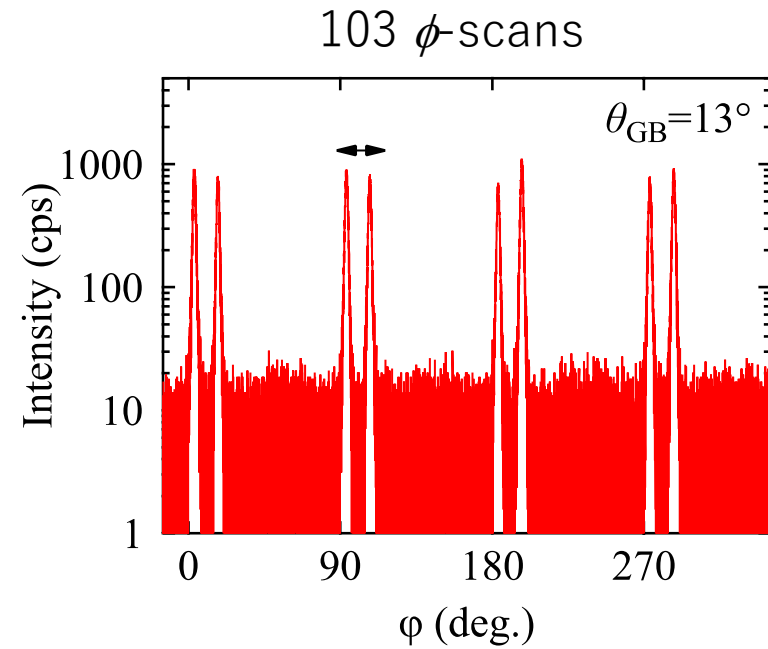
K-doped Ba122 grown on Ba122-buffered MgO(001)



- The 00/ ω peaks arising from K-doped Ba122 and Ba122 were observed
- K-doped Ba122 was grown epitaxially on Ba122-buffered MgO
- T_c of K-doped Ba122 was 37.5 K

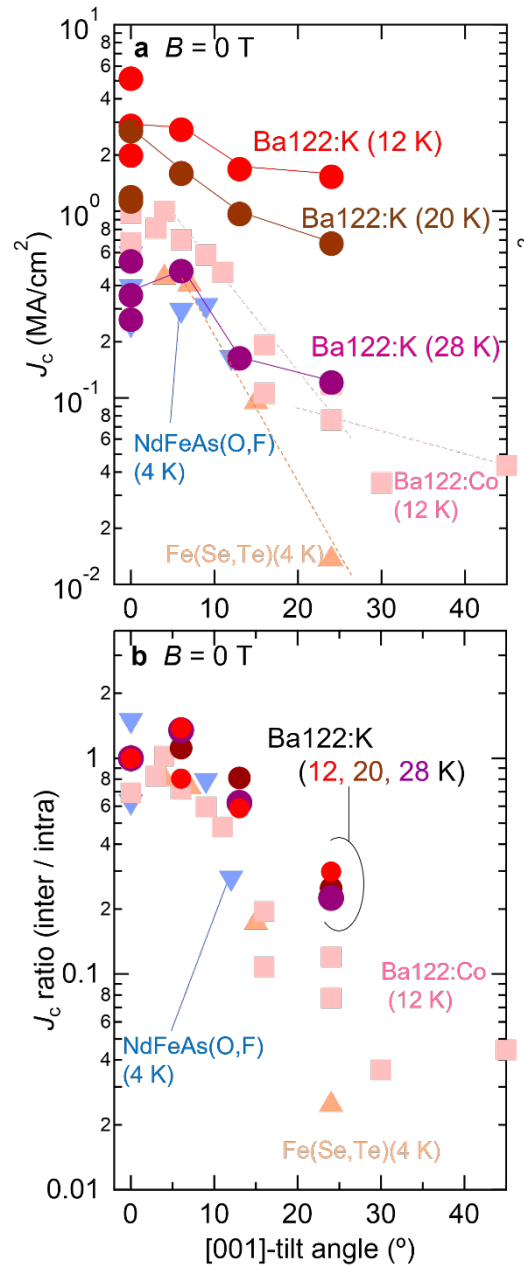
K-doped Ba122 bicrystals are realised

$$\theta_{GB} = \underline{13^\circ}$$



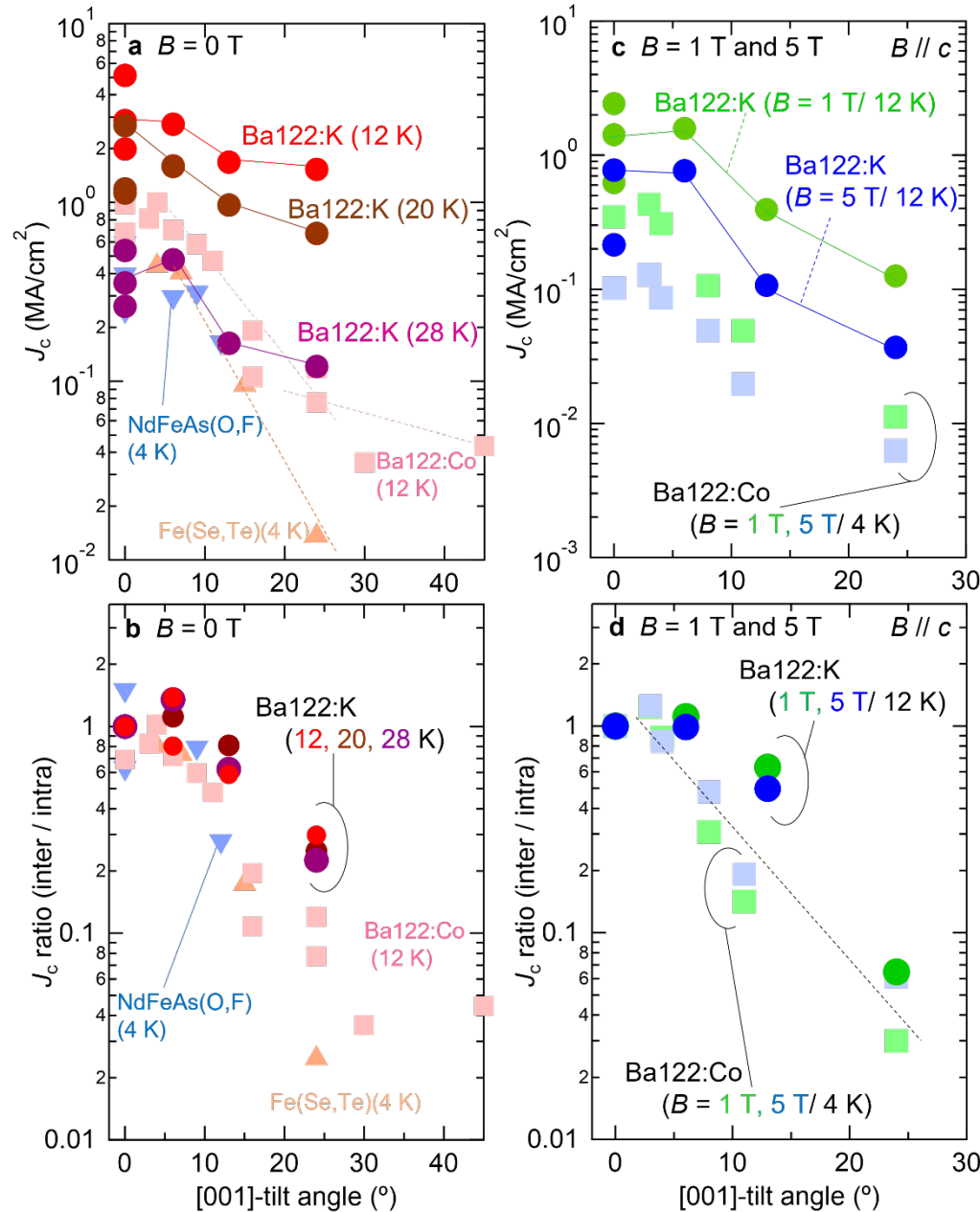
- Structurally fine K-doped Ba122 films are grown on MgO bi-crystal substrates
- Microbridge was fabricated by all dry processes (Ar ion etching or laser cutting)

Transport properties of K-doped Ba122 bicrystals



- Even at 12 K and 20 K, inter-grain J_c of K-doped Ba122 is higher than those of other IBS [Co-doped Ba122 (12 K), Fe(Se,Te) (4 K), NdFeAs(O,F) (4 K)]
- Critical grain boundary angle θ_c of (Ba,K)122 is $\sim 9^\circ$, similarly to other

Transport properties of K-doped Ba122 bicrystals



T. Hatano *et al.*, in preparation

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- Critical grain boundary angle θ_c of (Ba,K)122 is $\sim 9^\circ$, similarly to other IBS
- The θ_c of (Ba,K)122 is unchanged by applied magnetic fields (c.f. Co-doped Ba122)

1. Over-doped NdFeAsO showed a high J_c and low anisotropy
-> The strategy for over-doping method can be applicable
2. Low angle GBs and their networks work as strong pinning centers in K-doped Ba122
3. Inter-grain J_c exceed 1 MA/cm² even at $\theta_{GB}=24^\circ$ and 12 K
4. The critical angle θ_c for all IBSs seem to be $\sim 9^\circ$
5. The θ_c for K-doped Ba122 is unchanged by magnetic field