# High Magnetic Field Generated by Bulk MgB<sub>2</sub> Prepared by Spark Plasma Sintering

Kévin Berger, Michael Rudolf Koblischka, Bruno Douine, Jacques Noudem, Pierre Bernstein, Thomas Hauet and Jean Lévêque

Abstract—From the applications point of view, the advantage of low density given by MgB2 material must be taken into consideration and, the generation of strong magnetic flux densities using MgB2 should be investigated. In this contribution, we have studied the magnetic properties of samples processed by a fast Spark Plasma Sintering machine that is able to produce dense and high quality MgB2 samples. Experiments were carried out both on small pieces of the samples and on large-sized samples performing trapped field measurements in a Field-Cooling process. For temperatures between 10 K to 30 K, the results show a strong dependence of the magnetic behavior of the large-sized samples to the applied magnetic field sweep rate, while nothing particular appeared on the small samples. As the magnetic flux density produced by MgB2 bulks is directly linked to the potential of the applications, we report the field produced at the surface of one single MgB2 sample of 30 mm diameter and the field produced inside a stack of two MgB2 samples of 20 mm in diameter. A generation of magnetic flux density up to 4.8 T @ 10 K and 3.92 T @ 20 K inside the stack of the two MgB2 samples was observed. According to these values, thanks to their very low density and to their ease of manufacturing, MgB<sub>2</sub> bulks are promising materials for the applications of superconductors.

Index Terms—Bulk MgB<sub>2</sub>, Field-cooling process, Flux jumps, Strong magnetic field, Trapped magnetic field.

# I. INTRODUCTION

In the field of High Temperature Superconductors (HTS), some applications requires the fabrication of large, bulk samples in order to employ them for levitation devices, fault current limiters, non-contact bearings for liquid pumping, magnetic shielding screens, and motors. MgB<sub>2</sub> material represents currently the metallic superconductor with the highest transition temperature [1]. A very promising feature of

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this material is that one can achieve a high critical current density in polycrystalline samples without aligning the crystal orientation with a substrate, which enables a much cheaper production technology [2]. Besides, different methods can be used for densifying the material such as Spark Plasma Sintering (SPS) process [3]. Using this process, MgB2 bulk samples with 99% of relative density have been obtained [4]. Theoretically, bulk MgB<sub>2</sub> samples have very low density of 2.64 g/cm<sup>3</sup>, compared to (RE)BaCuO samples which are around 6.36 g/cm<sup>3</sup>, and NdFeB permanent magnets at around 7.60 g/cm<sup>3</sup>. Practically, this ratio higher than 2-3 is maintained. From the applications point of view, this advantage must be taken into consideration and, the generation of strong magnetic flux densities using MgB2 should be investigated, even though MgB<sub>2</sub> requires a lower operation temperature. Nowadays, with developed cryo-cooling systems, this is not a large obstacle anymore.

In this contribution, we studied the magnetic properties of samples processed at CRISMAT by a fast SPS machine that is able to produce dense and high quality MgB<sub>2</sub> samples. Experiments were carried out both on small pieces of the samples and on large-sized samples performing trapped field measurements in a Field Cooling (FC) process.

#### II. SAMPLES CHARACTERIZATION

Magnetic characterization of a small sample piece of  $1.377 \text{ mm} \times 1.473 \text{ mm} \times 0.638 \text{ mm}$  was carried out at IJL Nancy using a Magnetic Property Measurement System (MPMS®3) from Quantum Design.

The DC magnetic moment m was measured at 10 K, 20 K and 30 K, for an applied field parallel to the longest dimension of the sample. Half of the Major Hysteresis Loops (MHL) of the magnetic moment m are shown in Fig. 1. The field was cycled between +5 T and -5 T with various sweep rates: 1 mT/s, 4 mT/s and 20 mT/s, but in this range no influence of the sweep rate was observed. It can also be mentioned that no flux jumps phenomena have been seen, even at high sweep rates.

# A. Irreversibility field H<sub>irr</sub>

The irreversibility field determined at the point where the increasing and decreasing branches of the loop merge together is equal to  $4.98\,$  T,  $3.33\,$  T and  $1.50\,$  T, respectively at  $10\,$  K,  $20\,$  K and  $30\,$  K.

## B. Determination of $J_c(B)$

The critical current density  $J_c$  is calculated from the Major

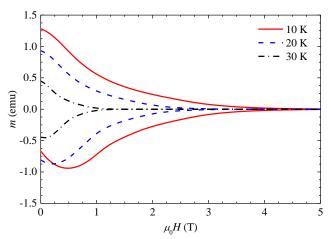


Fig. 1. Half of the Major Hysteresis Loops of the magnetic moment m of a small MgB<sub>2</sub> sample measured at 10 K, 20 K and 30 K, for an applied field cycled between +5 T and -5 T.

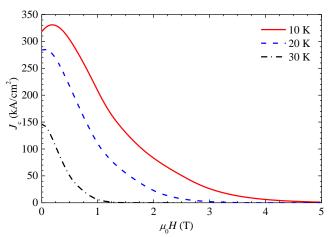


Fig. 2. Field dependence of the critical current density calculated at  $10~\rm{K}$ ,  $20~\rm{K}$  and  $30~\rm{K}$  from the MHL in Fig. 1.

Hysteresis Loop (MHL) for which the sample is in the critical state. In the case of an infinitely long sample with a rectangular cross-section  $(2a \times 2b)$ , with a < b, the critical current is given by

$$J_{c}(B) = \frac{\Delta M}{a(1 - a/3b)} \tag{1}$$

where  $\Delta M$  is the width of the curve [5], [6].

It can be noticed that due to some uncertainty in the size measurements of the sample and its volume estimation, the value of the critical current density are probably majored by 5%. Nevertheless, the critical current densities at 10 K and 20 K exhibit values higher than 10<sup>5</sup> A/cm<sup>2</sup> for 1 T of applied field, what is a relatively good result.

# III. EXPERIMENTAL SETUP

# A. Apparatus

The cryostat employed for the trapped field measurements in our case is an Oxford Instruments 5 T low-loss cryostat with a room-temperature bore. On removing the secondary cryostat for sample cooling used for low-temperature STM and MFM work, the room-temperature bore of 7.5 cm diameter is fully exposed. Fitting a cryocooler under it enables

trapped-field experiments to be performed in magnetic fields up to  $\pm 5$  T perpendicular to the sample surface. For our experiments, we have employed the ARS-4K-cryocooler system, which can reach temperatures down to 4.2 K but the lowest temperature in our experiments is roughly 9 K.

A bi-directional power supply for the superconducting coil (Oxford Instruments IPS 120-10) enables a continuous field sweep through 0 T together with a controlled magnetic field-sweep rate ranging up to 40 mT/s (2.4 T/min).

Up to 3 Hall probes may be employed to measure the fields at the sample surface. We have used high linearity Hall probes HHP-NP from Arepoc which have a sensibility higher than 70 mV/T. The overall dimensions of the Hall probe are 7 mm  $\times$  5 mm  $\times$  1 mm and the active area is approximately 0.35 mm far from the sensor surface. The calibrations of these sensors were performed at each mounting of a sample.

Fig. 3 (a) presents the sample stage of the cryocooler with two bulk MgB2 samples in place. In Fig. 3 (b), a copper ring is fitted around the sample for better temperature exchange during the cooling, and another copper plate is then fixed above the sample containing three slits for the Hall probes. Finally, Fig. 3(c) shows the entire cryocooler head with a sample and the Hall probes in place. In this way, samples up to 4 cm in diameter and 3 cm in length can be cooled and submitted to a high magnetic field.

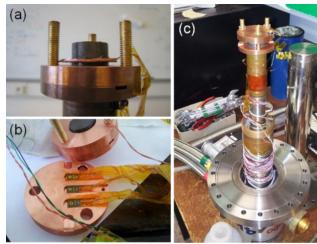


Fig. 3. (a) Head of the cryocooler with two bulk MgB<sub>2</sub> samples in place. (b) Copper top-plate with slits for three Hall probes spaced by 1 cm. (c) The entire cryocooler system with a mounted sample.

# B. Large-sized MgB2 samples

Three large-sized  $MgB_2$  samples have been used in our trapped field measurements. The height h of each of them is 9.7 mm and they have been made using a SPS process with a heat treatment at 1 200 °C under 50 MPa pressure.

One sample, called hereafter d30, with a diameter d of 30 mm, has been reacted during 15 min. The two others samples, named d20, with a diameter of 20 mm have been reacted during 10 min each.

# IV. TRAPPED FIELD MEASUREMENTS

In the following, we present the trapped field measurements on the large-sized bulk  $MgB_2$  samples.

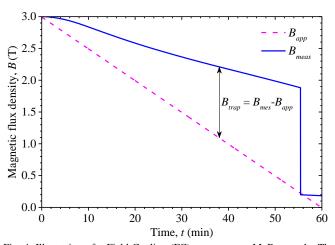


Fig. 4. Illustration of a Field Cooling (FC) process on a MgB<sub>2</sub> sample. The difference between the measured field  $B_{\rm meas}$  and the applied field  $B_{\rm app}$  is defined as the trapped magnetic field  $B_{\rm trap}$ .

# A. Field-cooling process

To activate/energize the MgB<sub>2</sub> pellets, we employ the Field-Cooling (FC) process. For this purpose, we tried also the use of different field sweep rates in order to see the dependence on potential flux jumps of this sweep rate. An example of FC process is illustrated in Fig. 4. Here, a magnetic field of 3 T is applied to the sample which is then cooled down below its critical temperature. The applied magnetic field is then decreased with a fixed sweep rate. The difference between the measured field  $B_{\text{meas}}$  and the applied field  $B_{\text{app}}$  is defined as the trapped magnetic field  $B_{\text{trap}}$ , this is the reaction of the sample. The remanent trapped magnetic field  $B_{\text{rem}}$  is defined when the applied field is null. The definition of the trapped magnetic field  $B_{\text{trap}}$  is essential, since only the remanent trapped magnetic field is usually reported. However for practical applications as in [7], higher values of  $B_{\text{trap}}$  are of great importance.

# B. One single MgB2 sample

In this section, we report the performances in terms of trapped magnetic field of one single MgB<sub>2</sub> sample.

Fig. 5 shows the magnetic field trapped at the center of the surface of the sample *d30* at 10 K and 20 K. At 20 K, there are

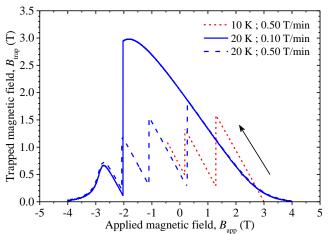


Fig. 5. Magnetic field trapped at the center of the surface of the sample d30 at 10 K and 20 K. At 20 K, there are less flux jumps when the field sweep rate is 0.10 T/min instead of 0.50 T/min.

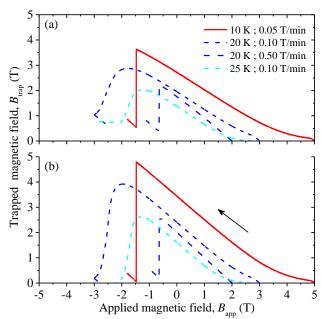


Fig. 6. Trapped magnetic field results on a stack of two MgB2 samples d20x2 at 10 K, 20 K and 25 K (a) at the center of the surface of the stack, (b) at the center between the two samples.

less flux jumps when the field sweep rate is 0.10 T/min instead of 0.50 T/min. This sweep rate influence on the flux jumps has also been observed in [8]. The remanent magnetic field is equal to 2.05 T at 20 K and the highest value of trapped/produced magnetic field by this sample is around 2.98 T at 20 K. Similar values around 2 T have been recently summarized for the same sample size and at the same temperature in Table 1 of [9].

Given that many flux jumps occurred at 10 K, even at low sweep rate, it was not possible to perform high trapped field at this temperature. More details concerning the flux jumps in MgB<sub>2</sub> can be found in [10].

## C. Two $MgB_2$ samples stacked together

In this section, we report the performances in terms of trapped magnetic field of two MgB<sub>2</sub> samples stacked together. Because we did not have two samples of 30 mm of diameter, we used samples of same properties but with a diameter of

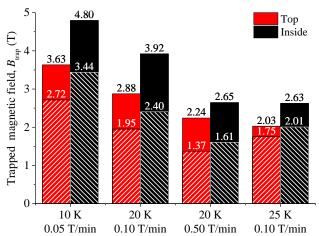


Fig. 7. Bar graph synthetising the results of trapped magnetic field using d20x2. The hashed part of the bars represents the remanent trapped magnetic field, whereas the full bar is associated to the maximum trapped field.

20 mm, that is the two d20 samples. The two d20 pellets are separated by an air gap of the size of the Hall probe, i.e. 1 mm. By neglecting this air gap, it can be considered that the stack forms a single MgB<sub>2</sub> pellet of 20 mm of diameter and roughly 20 mm height, that will be later called d20x2.

Fig. 7 shows a bar graph synthetising the results of trapped magnetic field using d20x2. The red bars (left) correspond to the measures at the center of the **top** surface of the d20x2 assembly, whereas the black ones (right) correspond to the measures **inside** d20x2, in the air gap. The corresponding values are directly written above the bars. The hashed part of the bars represents the remanent trapped magnetic field, when the applied field is null, whereas the full bar is associated to the maximum trapped field achieved in our experiments. Similar results on others samples have been reported in [11], [12].

## V. RESULTS AND DISCUSSION

It is obvious that the remanent magnetic field anywhere in the sample cannot be higher than the irreversibility field reported in Section II.A. The irreversibility field values of our samples prepared by SPS are clearly lower than those reported in [9] for samples prepared by hot-pressing of ball-milled precursor powder, where the irreversibily fields are about 7 T at 20 K and 2 T at 30 K.

According to [13]–[15], the external magnetic field necessary for full flux penetration at the center of the sample can be written as

$$B_{p} = \mu_{0} J_{c}^{\text{bulk}} R f(k_{d}), \text{ with } f(k_{d}) = k_{d} \ln \left( \frac{1 + \sqrt{1 + k_{d}^{2}}}{k_{d}} \right) (2)$$

with  $J_c^{\text{bulk}}$  as the field-independent bulk critical current density, R the radius, d the diameter and h the height of the sample, and  $k_d = h / d$ . As this field  $B_p$  corresponds to the full penetration field, it gives an idea of the achievable trapped magnetic flux density at the center of the sample without taking into account the  $J_c(B)$  law of the material, i.e. for a constant critical current density. Using the same hypotheses, one can also express the corresponding field at the center of the surface of the sample

$$B_{\rm p}^{\rm s} = \frac{\mu_0 J_{\rm c}^{\rm bulk} R}{2} f(k_R), \text{ with } k_R = 2k_d = h/R$$
 (3)

In that case, the ratio  $B_p / B_p^s$  is only driven by geometric parameters. For the samples studied here, this ratio is equal to 1.51, 1.63 and 1.82 for d30, d20 and d20x2, respectively. As  $\lim_{k\to\infty} f(k) = 1$ , a 'half' infinite long cylinder would correspond to a ratio  $B_p / B_p^s$  of 2. However, extrapolating the value of the trapped magnetic field at the center of the sample, as some authors do, by multiplying the value of the trapped magnetic field at the surface with the theoretical ratio  $B_p / B_p^s$  is not sufficient, and leads to overestimated values. The comparison is made in Table I. Given the fact that the only relevant value for the practical applications as permanent magnet is the trapped magnetic field at the surface of the sample, the value of the trapped field inside the sample is usually reported for

TABLE I COMPARISON OF THEORETICAL AND EXPERIMENTAL RATIO  $B_P / B_P^S$ 

Ratio $B_p / B_p^s$	10 K 0.05 T/min	20 K 0.10 T/min	20 K 0.50 T/min	25 K 0.10 T/min
Theoretical	1.63	1.63	1.63	1.63
With $B_{\text{trap}}^{\text{max}}$	1.32	1.36	1.18	1.30
With B <sub>rem</sub>	1.26	1.23	1.18	1.15

comparison reasons or in order to impress.

From (2) and (3), one can say that the trapped magnetic field, for a fixed height of the sample, can be increased simply by increasing the diameter of the sample. Theoretically, choosing d30 instead of d20 leads to 26% and 35% of increase for  $B_p$  and  $B_p{}^s$ , respectively. In the same way, choosing d20x2 instead of d20 leads to 23% and 10% of increase for  $B_p$  and  $B_p{}^s$ , respectively. Therefore, the increase of the trapped magnetic field at the surface of the sample is more significant by increasing the diameter from 20 mm to 30 mm that by doubling the height of d20. Finally, choosing d30 instead of d20x2 leads to 2% and 23% of increase for  $B_p$  and  $B_p{}^s$ , respectively. The same tendency is reported experimentally at 20 K, the increase of the maximal trapped field from d30 to d20x2 is about 3%.

From (3), the field-independent bulk critical current density  $J_c^{\text{bulk}}$  can be calculated. For d20x2 at 20 K, with 2.88 T of maximum trapped magnetic field at the surface of the sample,  $J_c^{\text{bulk}}$  is estimated around 48 kA/cm<sup>2</sup>. This value corresponds with the  $J_c(B)$  curve shown in Fig. 2 at 20 K and for B = 1.6 T. Therefore, the mean value of trapped magnetic field in the sample volume is approximately equal to this value of 1.6 T.

The values obtained here of the remanent trapped magnetic field measured at the center of the surface of the sample are a little bit lower than highest reported up to now in [9], i.e. 4.6 T, 3.3 T and 2 T, at 15 K, 20 K and 25 K, respectively. However, the obtained values of the maximal trapped magnetic field for d20x2 of 3.63 T at 10 K, 2.88 T at 20 K (2.98 T for d30) and 2.03 T at 25 K are really interesting for the intended applications.

#### VI. CONCLUSION

Through this article, we have deeply investigated the trapped magnetic field, not only at the surface of a sample produced by SPS, but also inside a stack of two  $MgB_2$  samples. The useful values of trapped magnetic field obtained at 20 K at the surface of the samples are really interesting for the intended applications. According to these values, to their very low density compared to other materials and to their ease of manufacturing,  $MgB_2$  bulks are promising materials for the applications of superconductors.

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